

# Sensitivity of Tropical Cyclones to Parameterized Convection in the NASA GEOS5 Model

**Young-Kwon Lim<sup>1,2</sup>, Siegfried D. Schubert<sup>1</sup>,  
Oreste Reale<sup>3,4</sup>, Myong-In Lee<sup>6</sup>, Andrea M. Molod<sup>1,5</sup>, Max J. Suarez<sup>1</sup>**

Corresponding author: Young-Kwon Lim (Young-Kwon.Lim@nasa.gov)

<sup>1</sup>Bldg. 33, code 610.1, 8800 Greenbelt Rd.,  
Global Modeling and Assimilation Office, NASA/GSFC, Greenbelt, Maryland, 20771

<sup>2</sup>Goddard Earth Sciences Technology and Research, I. M. System Group

<sup>3</sup>Earth Science Division, Atmospheres, NASA/GSFC, Greenbelt, Maryland

<sup>4</sup>Goddard Earth Sciences Technology and Research,

Universities Space Research Association (USRA)

<sup>5</sup>ESSIC, University of Maryland, College Park, Maryland

<sup>6</sup>Ulsan Institute of Science and Technology (UNIST), Ulsan, South Korea

February 03, 2014

Submitted to Journal of Climate

Abstract

The sensitivity of tropical cyclones (TCs) to changes in parameterized convection is investigated to improve the simulation of TCs in the North Atlantic. Specifically, the impact of reducing the influence of the Relaxed Arakawa-Schubert (RAS) scheme-based parameterized convection is explored using the Goddard Earth Observing System version5 (GEOS5) model at 0.25° horizontal resolution. The years 2005 and 2006 characterized by very active and inactive hurricane seasons, respectively, are selected for simulation.

A reduction in parameterized deep convection results in an increase in TC activity (e.g., TC number and longer life cycle) to more realistic levels compared to the baseline control configuration. The vertical and horizontal structure of the strongest simulated hurricane shows the maximum lower-level (850-950hPa) wind speed greater than 60 m/s and the minimum sea level pressure reaching ~940mb, corresponding to a category 4 hurricane - a category never achieved by the control configuration. The radius of the maximum wind of ~50km, the location of the warm core exceeding 10°C, and the horizontal compactness of the hurricane center are all quite realistic without any negatively affecting the atmospheric mean state.

This study reveals that an increase in the threshold of minimum entrainment suppresses parameterized deep convection by entraining more dry air into the typical plume. This leads to cooling and drying at the mid- to upper-troposphere, along with the positive latent heat flux and moistening in the lower-troposphere. The resulting increase in conditional instability provides an environment that is more conducive to TC vortex development and upward moisture flux convergence by dynamically resolved moist convection, thereby increasing TC activity.

## 1. Introduction

This article is inspired by a recent research on tropical cyclone (TC) simulation coordinated by the US CLIVAR hurricane working group (<http://www.usclivar.org/working-groups/hurricane>). Among various science issues raised in the research, it was found that many current general circulation models (GCMs) seriously underestimate tropical cyclone (TC) activity over the North Atlantic when run at  $\sim 0.5$  degree latitude/longitude or coarser horizontal resolution as opposed to some other basins such as the North Pacific (Walsh et al. 2013; Shaevitz et al. 2014). Only a few recent GCMs simulate TC activity in the North Atlantic with some success (e.g., Zhao et al. 2009). A question is raised on whether horizontal resolution finer than  $\sim 0.5$  degree is necessary to achieve reasonable TC numbers over the North Atlantic, even though resolution appears less critical for simulating TC numbers than the intensity (Strachan et al. 2013). Compared to TC numbers, simulating realistic TC intensity and spatial structure appears more unattainable at these resolutions ( $\sim 0.5$  degree or coarser), with generally weak and poorly organized systems occurring in all the GCMs (more than 10) involved in the working group. Higher resolution may be more appropriate to allow for a better representation of African easterly waves and the associated strong low-level vorticity, leading to more realistic simulation of Atlantic TC activity including number, intensity, life-cycle and spatial structure (Knutson et al. 2007; Caron et al. 2011; Putman and Suarez 2011; Daloz et al. 2012; Emanuel 2013).

Several previous studies have identified deficiencies in the simulation of moist convective processes as substantially affecting the ability to simulate TC characteristics (Slingo et al. 1994; Zhang and McFarlane 1995; Smith 2000; Benjamin et al. 2008; Zhao et al. 2012). Smith (2000) suggests that major improvements in TC predictions will depend on improvements in the

representation of convection in hurricane models. Knutson et al. (2004), Atlas et al. (2005), LaRow et al. (2008), Ma and Tan (2009) and Reed and Jablonowski (2011) performed model experiments to examine the sensitivity of the global TC numbers to changes in the convective parameterizations. Shen et al. (2006) and Stan (2012) examined the impact of disabling the convective parameterizations (in particular the schemes developed by Arakawa and Schubert (1974) and Moorthi and Suarez (1992)) on GCM hurricane forecasts. Both studies agreed on that explicit representation of cloud processes produces a larger number of TC events, with stronger intensity and longer life-cycles (Reed and Jablonowski 2011; Stan 2012). However, the details of the atmospheric processes responsible for altering TC activity were not the focus of the above studies. Some of the atmospheric responses to changes in deep convective activity are discussed in Zhao et al. (2012), which focuses on the impacts on vertical velocity, temperature and humidity. That study identifies stronger sensitivity of mid-tropospheric vertical motion to the minimum entrainment compared with the other variables (humidity and vertical wind shear): the importance of the response of vertical motion is also discussed in Oouchi et al. (2006).

The sensitivity of the atmospheric response to the convective parameterization and its impact on TC activity motivates the current study to improve the simulation of TC activity over the North Atlantic in GCMs. As in Tokioka et al. (1988), Held et al. (2007) and Kang et al. (2008), we find that increasing the threshold of minimum entrainment rate makes the convective plume entrain more dry air as it ascends, losing in-cloud buoyancy more rapidly when entraining. As a result, deep convection is inhibited from occurring in updrafts with a lateral entrainment rate  $\lambda$  lower than a threshold value  $\lambda_0$  (Tokioka et al. 1988; Held et al. 2007). Parameterized deep convection, including the sub-grid scale cloud representation is restricted by increasing  $\lambda_0$  (Held et al. 2007; Kang et al. 2008), while the fraction of large-scale condensation by non-

convective clouds increases through the large-scale cloud/condensation module (Tiedtke 1993) as opposed to the convection module (Kang et al. 2008). Wang (2014) found that this non-deep convection plays a role in moistening the lower to middle troposphere whereas deep convection moistens the upper-troposphere. Zhao et al. (2012) also found that an inhibition of the convective parameterization through enhanced lateral mixing into convective plumes leads to increased TC genesis, along with a colder and drier upper troposphere.

The main objective of this study is to improve the simulation of North Atlantic TC activity in the National Aeronautics and Space Administration (NASA) Goddard Earth Observing System Version 5 (GEOS-5) model (Rienecker et al. 2008; Molod et al. 2012), including the TC number, intensity, life cycle and horizontal/vertical structures. We focus on the cumulus entrainment rate in the GEOS-5 model run at a relatively high horizontal resolution (0.25 degree) to determine the thresholds of minimum entrainment that reliably reproduce the TC activity over the North Atlantic. An important aspect of this study is that we also focus on maintaining a realistic mean climate state – an aspect of the simulations (the response of the mean state to the changes in deep convection) many previous studies have not emphasized. The analysis of the runs is geared to better understanding the atmospheric response determining the TC activity over the North Atlantic with a particular focus on the atmospheric instability over the TC genesis region, the thermodynamic and radiative balance, the low-level fluxes and circulation, and the mean climate state response.

The organization of the paper is as follows. Section 2 describes the GEOS-5 model and the experimental design. The results of the sensitivity experiments are presented in Section 3, including TC numbers, intensity, life cycle and spatial structure of the strongest hurricanes. We investigate in Section 4 the relevant atmospheric dynamics and physics to explain the TC activity

changes that occur as a result of the modifications to the parameterized deep convection. Some verification of the atmospheric basic state changes due to the deep convection changes is also provided in Section 4, followed by concluding remarks and discussion in Section 5.

## 2. Model and Experimental design

### 2.1 Model

We utilize the NASA GEOS-5 model (Rienecker et al. 2008; Molod et al. 2012) for our experiments. The model is run with 72 hybrid-sigma vertical levels, extending to 0.01 hPa, and  $\sim 0.25^\circ$  latitude/longitude horizontal resolution. The convection scheme in the GEOS5 model is basically a modified version of the Relaxed Arakawa Schubert (RAS) scheme of Moorthi and Suarez (1992). In the modified version, multiple convective plumes in RAS have the convection base level at the top of the planetary boundary layer. The adjustment time scale of plumes is a function of the vertical depth of plumes, which varies from 0.5 hr for the shallowest plume to 12 hr for the deepest. Another major modification to RAS, which affects convective variability of the model substantially and hence the intensity of simulated TCs, is the stochastic determination of the minimum entrainment threshold. This is basically same as in Tokioka et al. (1988), which modification is known as limiting frequent adjustment of deep convection in a model and increasing the variability of deep convection (Lee et al. 2003; Lin et al. 2008). For “Tokioka” constraint in this study, the threshold value is non-deterministic so that the subsequent value is determined randomly with a predetermined power function:

$$\lambda_0 = \frac{\alpha}{D} x^n, \quad 0 \leq x \leq 1 \quad (1)$$

where  $\alpha$  is a free parameter,  $D$  the depth of the subcloud layer, and  $x$  a random number chosen at every model integration time step. Here the depth of the subcloud layer  $D$  is assumed to be

related with the plume radius by Simpson and Wiggert (1969), and this stochastic determination has a basis to consider spatial statistics of the observed cloud widths (or radius) and volumes that follow a power law distribution found by Bacmeister and Stephens (2011) from the CloudSat data (Stephens et al. 2002). Figure 1 compares the minimum entrainment threshold values between the conventional, deterministic method from Tokioka et al. (1988) and the stochastic method in this study. In case of  $n=1/2$ , which is used in this study, the threshold curve increases with the square root function (Fig. 1), and restores to the conventional Tokioka constraint. This tends to impose restriction on the entrainment rate in most cases, but occasionally allowing a very deep and less entraining convection plumes. Putman and Suarez (2011) used the same version of GEOS-5 that used in this study for their tropical cyclone simulation, but with much higher horizontal resolution of  $\sim 7\text{km}$ .

Increase in the Tokioka constraint (i.e., increase in the threshold of minimum entrainment) makes more difficult condition for sub-grid scale deep convection. While the parameterized deep convective process is suppressed as a result of this Tokioka constraint, resolved-scale non-convective clouds are favored and thus, more active large-scale cloud/condensation and precipitation process becomes feasible.

## 2.2 Experimental design

Three types of experiments are carried out by applying different Tokioka constraints in model simulations of the years 2005 and 2006, which are very active and inactive hurricane years, respectively. We call these three experiments A, B, and C. Experiment A (ExpA) has the strongest Tokioka constraint, where the greatest possibility of large thresholds of minimum entrainment exists, while Experiment B corresponds to an intermediate range of thresholds

(ExpB), and Experiment C (ExpC), which is a baseline control run, corresponds to the lowest minimum entrainment range that allows more parameterized deep convection relative to the ExpA and ExpB. These were implemented in the model with the minimum  $\alpha/D = 2 \text{ km}^{-1}$  (ExpA),  $0.7 \text{ km}^{-1}$  (ExpB), and  $0.45 \text{ km}^{-1}$  (ExpC), respectively in Equation (1). Please note that solid curve as a bottom line for minimum entrainments in Fig. 1 can vary with the  $\alpha/D$  values given in each experiment. The curve for ExpA is placed above the other two curves (for ExpB and ExpC), while the opposite is true for the curve for ExpC. Although the choice of  $n$  in Equation (1) can affect the results, but this is left for future sensitivity study. As mentioned earlier, the horizontal resolution is a quarter-degree latitude/longitude. The minimum entrainment thresholds is increased in our experiments (e.g., ExpA and ExpB) to identify the model sensitivity in simulating TCs with a focus on better performance across the Atlantic basin. From the above experiments, we intend to estimate the best Tokioka constraint as determined by the improvements in TC simulation including genesis, lifetime, track and intensity at this resolution. Each type of experiments consists of three ensemble members. All simulations are done with prescribed weekly sea surface temperature (SST) forcing (the HadISST of Rayner et al. 2003), and are initialized from Modern-Era Retrospective Analysis for Research and Applications (MERRA, Rienecker et al. 2011) data on May 1 each year.

### 3. Sensitivity of TC characteristics to changes in parameterized deep convection

TCs are detected for each year (2005 and 2006) and each experiment (ExpA, ExpB and ExpC). The number of detected tropical storms and hurricanes is shown in Table 1. The basic algorithm of TC detection and tracking is the same as that used in Vitart et al. (2003), Knutson et al. (2007) and LaRow et al. (2008). Threshold values for wind speed, relative vorticity, warm



core, sea level pressure and minimum duration time are assigned in the algorithm to define TC (Table 2) at a quarter degree resolution based on suggestions in Walsh et al. (2007). The threshold values used here can be considered realistic because they are very close to the thresholds to detect TCs from observational data. ExpA produces an excessive number of tropical storms and hurricanes, compared with observations shown on the rightmost column (Table 1). In particular, it produces too strong TC activity (hurricanes more than 20) even in 2006, which was recorded as a relatively calm hurricane year (Knapp et al. 2010). ExpB produces more reasonable TC numbers, as they are closer to observations than either ExpA or ExpC. ExpC shows generally weak TC activity and produces much weaker than observed TC activity during the active hurricane season of 2005.

TC tracks are plotted for each experiment. During the 2005 season (Fig. 2), a variety of tracks were observed: 1) Cap Verde-type TC systems that crossed the Atlantic making landfall over the American continent, 2) early recurvers, 3) Gulf of Mexico TCs, and also 4) eastward tracking systems embedded in the westerlies. Comparison between the experiments indicates that ExpA and ExpB both display a variety of tracks much larger than ExpC, and closer to the observation. The results for 2006 (Fig. 3) show features similar to those obtained for 2005 in the sense that the increased thresholds of minimum entrainment produce larger number of TCs with longer life cycles and a track variety more similar to the observation over the Atlantic. This model response is consistent with Emanuel et al. (2008) and Stan (2012), that showed a sensitivity of TC activity to changes in the cumulus parameterization in the Community Climate System Model (CCSM) (Gent et al. 2011).

Additional details of the model responses (including TC activity, TC intensity and vertical hurricane structure) produced in each experiment are shown in Fig. 4 through 8. Scatter plots of

low-level (850-950hPa) wind versus SLP for 2005 (Fig. 4), show that the maximum wind speed is up to  $61\text{ms}^{-1}$  in the ExpA,  $60\text{ms}^{-1}$  in ExpB, and less than  $60\text{ms}^{-1}$  in ExpC. The minimum SLP is down to 936, 943, and 963hPa, respectively in the three experiments. ExpA and ExpB reveal improved distributions of SLP and wind speed compared with ExpC. Note however that the observed wind speed is up to  $80\text{ms}^{-1}$  and minimum SLP deeper than 920hPa for extremely strong hurricanes – something not reproduced by the model in any of the experiments. Much finer horizontal resolution (finer than 10km) such as that of Noda et al. (2012) and Putman and Suarez (2011) appear to be necessary for capturing such extremely strong hurricanes. The maximum wind speeds identified in ExpA and ExpB are comparable to those found in the idealized hurricane experiment of Reed and Jablonowski (2011). They simulated maximum wind speeds of  $\sim 60\text{ms}^{-1}$  at  $\sim 1\text{km}$  height and a quarter-degree horizontal resolution.

Figure 5 is the same as Fig. 4 but for 2006. The results show that the simulated wind speeds and minimum center SLP deepening for very strong hurricanes is comparable to or stronger than those for the observations (ExpA and ExpB). In contrast, ExpC exhibits weaker maximum wind and minimum SLP distribution compared with the observations.

The vertical structure of the strongest hurricane for ExpA is plotted in Fig. 6. The top left panel shows the wind speed (shaded) and temperature (contoured), and the right panel shows the warm core. The results for 2005 show a compact core and hurricane eye that is well defined. The maximum low level wind speeds (greater than  $60\text{ms}^{-1}$ ) and the radius of maximum winds (less than 50km) both correspond to very realistic TC representation. The warm core value is greater than  $10^{\circ}\text{C}$  and situated in the upper-troposphere, quite similar to the typical structure of observed hurricanes (Frank 1977). The bottom panel is the time series of the minimum SLP and maximum wind speed following the moving hurricane. The minimum SLP drops to 935hPa and maximum

wind speed reaches  $61\text{ms}^{-1}$ . The main features for 2006 are quite similar to those found for 2005, showing a well-defined hurricane characterized by a maximum wind speed of  $65\text{ms}^{-1}$  and minimum SLP of 935hPa.

The vertical structures of the strongest hurricanes in ExpB are shown in Fig. 7. Compared with ExpA, they are a little weaker, but still show a reasonable vertical structure in terms of compactness of the core, with the radius of the maximum wind less than  $\sim 50\text{km}$  (2 grid points in longitudinal direction), the maximum wind value near  $60\text{ms}^{-1}$ , the minimum SLP of about 940hPa, and vertical profile of the warm core (reaching  $10^\circ\text{C}$ ), which corresponds to a category 4 hurricane based on the Saffir-Simpson scale (<http://www.nhc.noaa.gov/aboutsshws.php>) in 2005. Observed wind .vs. SLP scatter plots in Figs. 4 and 5 indicated that the strongest observed hurricane in 2006 is weaker than that in 2005. The model reproduces well this difference between the two years (Fig. 7), but the difference is not as distinguishable as that found in the observation. As discussed in Fig. 4, the observed category 5 hurricane in 2005 was very hard to reproduce through a quarter-degree resolution run, resulting in rather a small difference in the intensity of the strongest hurricane between the two years in Fig. 7.

ExpC, which is the control run, also shows reasonable results (Fig. 8), though the storms are less organized than those found in the other two experiments. The warm core is weak with a value of approximately  $6\sim 8^\circ\text{C}$  and located in the mid-troposphere. The wind speed, SLP, and sharpness of the hurricane core are also not as realistic as those in ExpA and ExpB. The weak hurricane structures found in ExpC are not inconsistent with the study of Vitart et al. (2001), which found that the vertical structures of hurricanes simulated by RAS tend to be less intense and have lower warm cores than those produced with moist convective adjustment (MCA) schemes (Manabe 1969), though the RAS scheme produces more TCs and higher relative

humidity than MCA schemes. In any event, the hurricane vertical structures shown in Fig. 6 through 8 provides compelling evidence that the suppression of parameterized deep convection tends to produce more intense hurricanes with better organized vertical structures.

#### 4. Atmospheric response to parameterized deep convection

The physical mechanisms by which the suppression of the cumulus parameterization promotes TC activity are investigated next. The role of atmospheric moist static stability – a factor known to affect TC characteristics (Smith 2000) is shown in Fig. 9, where the vertical profile of moist static energy (left panel) and the vertical gradient of equivalent potential temperature (right panel) are computed over the TC genesis region, which covers most of the central and eastern tropical Atlantic (60°W~15°W, 5°N~20°N). Profiles from ExpA, ExpB, and ExpC all exhibit a decrease with height in moist static energy at low levels (below 700~800hPa height) and an increase at upper-levels, implying the possibility of conditional instability. The ExpA result shows the possibility of stronger instability than ExpB and ExpC. The vertical gradient of equivalent potential temperature ( $-\frac{\partial \theta_e}{\partial p}$ ) (Fig. 9b) also shows that lower tropospheric vertical column with the negative  $-\frac{\partial \theta_e}{\partial p}$  is more extensive in Exp A than the other two experiments, indicating the most unstable atmosphere occurs in ExpA, followed by ExpB and ExpC.

The vertical temperature and humidity profiles are plotted over the TC genesis region in Fig. 10 to indicate how they are linked to the atmospheric instability. Here we plot the differences in the vertical profiles of temperature/humidity between two experiments. The most striking feature of these results is the mid- to upper-tropospheric cooling (Fig. 10a) and drying (Fig. 10b) and lower-tropospheric moistening (Fig. 10b) resulting from the reduction in parameterized deep

convection. This vertical structure of the humidity is in good agreement with the results of Zhao et al. (2012). With a larger threshold of minimum entrainment, it is not unexpected that more dry air will be entrained into the plume, resulting in a drier upper atmosphere with decreased buoyancy (Held et al. 2007). The reduction of in-cloud buoyancy leads to the suppression of parameterized deep convection and less upper-tropospheric condensation heating. Figure 10 reflects such a change in the vertical profile of temperature and moisture, providing favorable conditions for developing an unstable atmosphere. Upward moisture flux and dynamically resolved convection are more likely to occur after the lower troposphere is moistened enough (Wang 2014) under this unstable atmospheric condition.

Figure 11 shows the latent heating and evaporative flux with a focus on the near surface moisture flux. We note that the differences between ExpB and ExpC exhibit very similar patterns though with smaller magnitude (Figure not shown). Figure 11a shows that the latent heat flux is significantly increased in ExpA throughout the Atlantic basin. The enhanced evaporative flux from the ocean (Fig. 11b) gives rise to the increase in latent heat flux. We will see later (Fig. 12) that the upward flux at the surface is also enhanced by increased ascending motion in the unstable atmosphere. This is consistent with Zhao et al. (2012), who found that surface heat fluxes including latent heat flux were influenced by vertical motion, and that this was important for making the resolved-scale convection more favorable for tropical cyclone genesis.

Figure 11c quantifies the increase in conditional instability we inferred from Fig. 9 through 11b, in terms of the convective available potential energy (CAPE). It shows that CAPE is significantly enhanced in ExpA (compared with the control) through the constraint on the parameterized deep convection. CAPE exists within the conditionally unstable layer of the troposphere, when a lifted parcel of air is warmer than the ambient air. It can be interpreted,

based on Fig. 10, that strong upper-tropospheric cooling overlying relatively warmer air at lower-levels accompanied sufficient moisture, contributes to a build up of CAPE. The larger CAPE in ExpA compared with the control run (ExpC) also demonstrates that dynamically resolving convective processes would be enhanced with the suppression of parameterized deep convection, resulting in an increase in large-scale condensation and precipitation (see the fractional precipitation change in Fig. 15).

Figure 12a clearly illustrates the strong response of the vertical motion (see also Oouchi et al. 2006) to the unstable atmosphere caused by the restriction of parameterized deep convection. The increase in atmospheric instability over the TC genesis region leads to increased ascending motion throughout the vertical column with the largest increase occurring in ExpA, followed by ExpB and then ExpC. The enhanced vertical ascent is associated with enhanced low-level convergence (Fig. 13). This low-level convergence, combined with low-level moistening (Fig. 10) and positive latent heat flux (Fig. 11a), plays a role in increasing the moisture flux convergence at low levels (Fig. 12b), with the largest values occurring in ExpA, followed by ExpB and ExpC.

Figure 13 shows the distribution of the differences between ExpA and ExpC in some other atmospheric quantities. The results show that low-level vorticity (850hPa), near-surface wind variance (10m), low-level convergence (925hPa) and vertically integrated moisture flux convergence are all enhanced over the eastern Atlantic. The enhancement is more pronounced over the TC genesis region between about 10°N and 15°N (for genesis latitudes, see also Figs. 2 and 3). These spatial patterns emphasize once again that, as the parameterized deep convection is suppressed through an increase in the minimum entrainment threshold, the resulting increase in

atmospheric instability makes dynamically resolved moist convection more active by enhancing low-level atmospheric motions (e.g., convergence, vorticity and wind variance).

To quantify that the atmospheric structures discussed so far provide a favorable TC genesis environment, we next calculate the genesis potential index (GPI). Here, the original version of GPI (Emanuel and Nolan 2004) has been modified following Murakami et al. (2011) to include the influence of vertical motion in contributing to TC genesis. The modified GPI<sup>1</sup> is given by

$$GPI = |10^5 \eta|^{3/2} \left( \frac{RH}{50} \right)^3 \left( \frac{V_{pot}}{70} \right)^3 (1 + 0.1 V_s)^{-2} \left( \frac{-\omega + 0.1}{0.1} \right), \quad (2)$$

where  $\eta$  is the absolute vorticity at 850hPa level,  $RH$  is the relative humidity in percent at 850hPa,  $V_{pot}$  is the maximum potential intensity (MPI;  $m s^{-1}$ ) defined in Emanuel (1995),  $V_s$  is the vertical wind shear ( $V(850hPa)$  minus  $V(200hPa)$ ), and  $\omega$  is the vertical wind velocity ( $Pa s^{-1}$ ) at 500hPa.  $V_{pot}$  is modified by Bister and Emanuel (1998) from the original version of Emanuel (1995) so that

$$V_{pot}^2 = \frac{C_k T_s}{C_d T_0} (CAPE^* - CAPE^b), \quad (3)$$

where  $C_k$  and  $C_d$  are the exchange coefficient for enthalpy and the drag coefficient, respectively.  $T_s$  is SST and  $T_0$  is the mean outflow temperature. Two CAPE-related quantities,  $CAPE^*$  and  $CAPE^b$  are the CAPE of the air displaced upward from saturation at sea level with reference to ambient air and the CAPE of the air at boundary layer, respectively.

The GPI formulation considers the influence of all the horizontal and vertical atmospheric structures and atmospheric instability discussed in reference to Fig. 9 through 13. As shown in Fig. 14, the GPI exhibits larger values in ExpA and ExpB than in the control run (ExpC). Large

---

<sup>1</sup> The source code for the FORTRAN version of the GPI calculation is available from [ftp://texmex.mit.edu/pub/emanuel/TCMAX/pcmin\\_revised.f](ftp://texmex.mit.edu/pub/emanuel/TCMAX/pcmin_revised.f).

values occur over the Atlantic in low-latitudes, the Caribbean Sea, and Gulf of Mexico, where a majority of TC genesis takes place.

An important consideration is whether the above modifications in the convective parameterization produce changes in the atmospheric mean state and its variability (e.g., seasonal cycle) that are within acceptable ranges. Slingo et al. (1994) and Zhang and McFarlane (1995) found some sensitivity of the tropical mean climate to the parameterized moist convection. Figure 15 shows the differences between each experiment in the mean precipitation for 2005. We plot both the large-scale precipitation and convective precipitation to indentify the fractional precipitation changes. The left panels show the difference between ExpA and ExpC in total precipitation (Fig. 15a), large-scale precipitation (Fig. 15b) and convective precipitation (Fig. 15c). The results show that both total and large-scale precipitation increase noticeably over the tropics for ExpA, whereas there is a substantial reduction in convective precipitation. In contrast, for ExpB (the right panels) there is only a very small change in total precipitation over the tropics (Fig. 15d), while the large-scale precipitation over the TC genesis region increases slightly (Fig. 15e) along with a decrease in convective precipitation (Fig. 15f). This fractional precipitation change is due to an increased activity of grid-scale cloud/condensation (i.e., large-scale precipitation) along with a suppression of parameterized deep convection (i.e., convective precipitation) (Reed and Jablonowski 2011; Zhao et al. 2012).

The area-averaged difference in total precipitation over the Atlantic basin between ExpB and ExpC (Table 3) is  $0.03 \text{ mm day}^{-1}$  in 2005 and  $0.09 \text{ mm day}^{-1}$  in 2006, while for ExpA the precipitation increases with respect to ExpC by 0.5 and  $0.6 \text{ mm day}^{-1}$  in those two years, respectively. All three experiments overestimate the observed precipitation (left column in Table 3; see also Fig. 16a). Since the observationally-based precipitation estimates from the Global



Precipitation Climatology Project (GPCP) version 2.2 (Adler et al. 2003; Huffman et al. 2009) are provided on a relatively coarse grid (2.5° lon.-lat.), we speculate that part of the apparent overestimation by the model might in fact represent an underestimation of the coarser resolution GPCP precipitation. Table 3 indicates that the excessive total precipitation in ExpA makes it a less than ideal model to use even though it has more intense hurricanes. Also, ExpA produces an excessive number of storms during the inactive hurricane year of 2006 (Table 1). In contrast, ExpB seems a better option since it has minimal impact on the total precipitation (amount and spatial distribution), and it has improved 3-dimensional hurricane structures, intensities, and numbers, compared with the control (ExpC).

Figure 16 compares for each experiment the seasonal variation (June through November) of precipitation, SLP, surface temperature and upper-level wind speed over the typical TC genesis region. We see that ExpB tends to remain close to ExpC in all the variables, while ExpA shows a noticeable increase in total precipitation and surface temperature, and decrease in SLP compared to the control run. A comparison with observations indicates that, except for the upper-level wind speed, ExpA tends to have larger biases than either ExpB or the control (ExpC). Overall, our results suggests that a small reduction in the influence of the RAS convective parameterization (ExpB) produces improved North Atlantic TC activity without having a negative impact on the mean climate of the model when run at ¼ degree spatial resolution.

## 5. Concluding remarks and discussion

This study investigated the sensitivity of the North Atlantic TC activity to changes in parameterized deep convection in the NASA GEOS-5 model run at 0.25 degree latitude/longitude horizontal resolution. The study found that reduction in the influence of the

RAS scheme resulting from an increase in minimum entrainment for parameterized deep convection improves/enhances (compared to the standard control model) the TC activity in terms of numbers, intensity, life cycle, and the 3-dimensional storm structures. Our case studies for the 2005 (very active) and 2006 (very inactive) hurricane seasons revealed that only a modest increase in minimum entrainment (ExpB) was necessary at a quarter-degree resolution to produce TC numbers reasonably close to observations. The strongest hurricane reached a minimum SLP of 940hPa and low-level maximum wind speeds greater than 60m/s. The vertical structure is characterized by a well-defined hurricane eye, an upper-tropospheric warm core value exceeding 10°C, and maximum winds located at low levels within 50km of hurricane center. These magnitudes are comparable to, or improvement over, those of Shen et al. (2006) and Reed and Jablonowski (2011) at similar resolution. Furthermore, we found that a stronger constraint on the parameterized deep convection (i.e., ExpA) produced even stronger hurricanes (with minimum SLP of 935hPa, and a maximum wind speed of 65m/s), but this strong reduction in parameterized convection was not desirable as it produced strong TC activity even in a weak hurricane year (2006) and lead to increased bias in the atmospheric mean state.

Our analysis of the causes of the TC changes suggests that the following processes are at work. First, an increase in the threshold of minimum entrainment causes more entrainment of dry air into the convective plume. This entraining process leads to a reduction in buoyancy, resulting in less favorable conditions for deep convection. This leads to a reduction in condensation heating in the upper-troposphere, resulting in upper-tropospheric cooling. As a result, the upper-troposphere becomes drier and cooler, while the lower-troposphere becomes moister and warmer. This is accompanied by an increase in positive near-surface latent heat flux. This change in the vertical structure of the atmosphere is apparent over the Atlantic basin, easing the condition for

developing atmospheric instability. The greater instability is evidenced by changes in the profile of moist static energy, the vertical gradient of equivalent potential temperature, and the distribution of CAPE. As a consequence to the changes in atmospheric instability there is an increase in a) low level positive vorticity and convergence, b) upward motion, and 3) low-level moisture flux convergence. These changes help the development of explicit-scale convection and subsequent large-scale condensation of wave vapor, and represent a transition of the moist deep convection from the parameterized kind to that occurring at the resolved scales. This is reflected by a decrease in convective precipitation contributed by parameterized convection and an increase in large-scale precipitation. The above interpretation is supported by the changes in the GPI (Emanuel and Nolan 2004; Murakami et al. 2011) showing (for ExpB) an enhancement in the TC activity over the Atlantic, with more and stronger TCs than the control run (e.g., category 4 hurricanes were produced, whereas the control never produced hurricanes stronger than category 3).

The explicit-scale moist convection favored by atmospheric instability appears to compensate for the upper-tropospheric cooling and drying produced by the reduction in parameterized deep convection, with the atmosphere acting to maintain thermal equilibrium in the upper-troposphere. This likely explains why, at least for the experiment with only a modest change in the threshold of minimum entrainment (ExpB), the total precipitation remains close to those of the control.

Due to the limited number of cases investigated here, it is premature to say that ExpB truly represents an optimal setting for the reliable simulation of TCs at this resolution. Nevertheless, the current set of experiments are encouraging in that they demonstrate how years with widely differing TC activity (2005 and 2006) can be reproduced by global climate model with a rather

minor change to the convective parameterization, and that this can be done without degrading the mean state.

This article suggests that quarter-degree horizontal resolution may be the minimum resolution necessary for achieving realistic TC simulations via modifications to the parameterized convective process (Reed and Jablonowski 2011), while still maintaining a realistic mean state. This conclusion is based in part on other experiments we carried out at half and one-third degree resolution (not shown), in which realistic TC numbers were difficult to achieve without modifying the threshold values for TC detection. Of course the results at  $\frac{1}{4}$  degree resolution still have several limitations. Key among them is the inability to produce the most intense (category 5) hurricanes. It appears that even at substantially finer than a quarter-degree resolution, fine scale hurricane structure can be improved through the modification of the convective parameterization. For example, Putman and Suarez (2011) concluded from the analysis of GEOS-5 forecasts run at 7-km horizontal resolution, that intense precipitation was better formed within the hurricane eye-wall and surrounding rain bands with a reduction in the influence of the RAS convective parameterization, compared to a control in which RAS was not modified. As such it appears that, even at considerably higher resolution than considered here, improvements in the simulation of TC activity can still benefit from a proper modification of the convective parameterization.

## References

- Adler, R. F., and Coauthors, 2003: The version 2 Global Precipitation Climatology Project (GPCP) monthly precipitation analysis (1979-present). *J. Hydrometeor.*, **4**, 1147-1167.
- Arakawa, A., and W. H. Schubert, 1974: Interaction of cumulus cloud ensemble with the large-scale environment. Part I. *J. Atmos. Sci.*, **31**, 671-701.
- Atlas, R., O. Reale, B.-W. Shen, S.-J. Lin, J.-D. Chern, W. Putmann, T. Lee, K.-S. Yeh, M. Bosilovich, and J. Radakovich, 2005: Hurricane forecasting with the high-resolution NASA finite volume general circulation model. *Geophys. Res. Lett.*, **32**, L03807, doi:10.1029/2004GL021513.
- Bacmeister, J. T., and G. L. Stephens, 2011: Spatial statistics of likely convective clouds in CloudSat data. *J. Geophys. Res.*, **116**, D04104, doi:10.1029/2010JD014444.
- Benjamin, M. S., C. Piani, W. Ingram, D. Stone, and M. Allen, 2008: Towards constraining climate sensitivity by linear analysis of feedback patterns in thousands of perturbed physics GCM simulations. *Clim. Dyn.*, **30**, 175-190.
- Bister, M., and K. A. Emanuel, 1998: Dissipative heating and hurricane activity. *Meteor. Atmos. Phys.*, **52**, 233-240.
- Caron, L.-P., C. G. Jones, and K. Winger, 2011: Impact of resolution and downscaling technique in simulating recent Atlantic tropical cyclone activity. *Clim. Dyn.*, **37**, 869-892.
- Daloz, A.-S., F. Chauvin, K. Walsh, S. Lavender, D. Abbs, and F. Roux, 2012: The ability of GCMs to simulate tropical cyclones and their precursors over the North Atlantic main development region. *Clim. Dyn.*, **39**, 1559-1576.

473 Emanuel, K. A., 1995: Sensitivity of tropical cyclones to surface exchange coefficients and a  
 474 revised steady-state model incorporating eye dynamics. *J. Atmos. Sci.*, **52**, 3969-3976.

475 Emanuel, K. A., and D. S. Nolan, 2004: Tropical cyclone activity and global climate.  
 476 Proceedings, *The 26th Conference on Hurricanes and Tropical Meteorology*, Miami, FL,  
 477 Amer. Meteor. Sci., 240-241.

478 Emanuel, K. A., R. Sundararajan, and J. Williams, 2008: Hurricanes and global warming:  
 479 Results from downscaling IPCC AR4 simulations. *Bull. Amer. Meteor. Soc.*, **89**, 347-367.

480 Emanuel, K. A., 2013: Downscaling CMIP5 climate models shows increased tropical cyclone  
 481 activity over the 21<sup>st</sup> century. *Proc. Nat. Acad. Sci.*, **110**, doi: 10.1073/pnas.1301293110.

482 Frank, W. M., 1977: The structure and energetic of the tropical cyclone. I. Storm structure. *Mon.*  
 483 *Wea. Rev.*, **105**, 1119-1135.

484 Gent, P. R., and co-authors, 2011: The Community Climate System Model Version 4. *J. Climate*,  
 485 **24**, 4973-4991.

486 Held, I. M., M. Zhao, and B. Wyman, 2007: Dynamic radiative–convective equilibria using  
 487 GCM column physics. *J. Atmos. Sci.*, **64**, 228–238.

488 Huffman, G. J., R. F. Adler, D. T. Bolvin, and G. Gu, 2009: Improving the Global Precipitation  
 489 record: GPCP version 2.1. *Geophys. Res. Lett.*, **36**, L17808, doi:10.1029/2009GL040000.

490 Kang, S. M., I. M. Held, D. M. W. Frierson, and M. Zhao, 2008: The response of the ITCZ to  
 491 extratropical thermal forcing: Idealized slab-ocean experiments with a GCM. *J. Climate*, **21**,  
 492 3521–3532.

493 Knapp, K. R., M. C. Kruk, D. H. Levinson, H. J. Diamond, and C. J. Neumann, 2010: The  
 494 international best track archive for climate stewardship (IBTrACS) : Unifying tropical

495 cyclone best track data. *Bull. Amer. Meteor. Soc.*, **91**, 363-376,  
 496 doi:10.1175/2009BAMS2755.

497 Knutson, T. R., and R. E. Tuleya, 2004: Impact of CO<sub>2</sub>-induced warming on simulated hurricane  
 498 intensity and precipitation: Sensitivity to the choice of climate model and convective  
 499 parameterization. *J. Climate*, **17**, 3477–3495.

500 Knutson, T. R., J. J. Sirutis, S. T. Garner, I. M. Held, and R. E. Tuleya, 2007: Simulation of the  
 501 recent multidecadal increase of Atlantic hurricane activity using an 18-km-grid regional  
 502 model. *Bull. Amer. Meteor. Soc.*, **88**, 1549–1565.

503 LaRow, T., Y.-K. Lim, D. Shin, E. Chassignet, and S. Cocke, 2008: Atlantic basin seasonal  
 504 hurricane simulations. *J. Climate*, **21**, 3191-3206.

505 Lee, M.-I., I.-S. Kang, and B. E. Mapes, 2003: Impacts of cumulus convection parameterization  
 506 on aqua-planet AGCM simulations of tropical intraseasonal variability. *J. Meteor. Soc.  
 507 Japan*, **81**, 963-992.

508 Lin, J.-L., M.-I. Lee, D. Kim, I.-S. Kang, and D. M. W. Frierson, 2008: Impacts of convective  
 509 parameterization and moisture convective trigger on AGCM-simulated convectively coupled  
 510 equatorial waves. *J. Climate*, **21**, 883-909.

511 Ma, L.-M. and Z.-M. Tan, 2009: Improving the behavior of the cumulus parameterization for  
 512 tropical cyclone prediction: Convection trigger. *Atmos. Res.*, **92**, 190-211.

513 Manabe, S., 1969: Climate and the ocean circulation: I. The atmospheric circulation and the  
 514 hydrology of the Earth's surface. *Mon. Wea. Rev.*, **97**, 739-774.

515 Molod, A., L. Takacs, M. Suarez, J. Bacmeister, I.-S. Song, and A. Eichmann, 2012: The GEOS-  
 516 5 Atmospheric General Circulation Model: Mean Climate and Development from MERRA

to Fortuna. *NASA Technical Report Series on Global Modeling and Data Assimilation*,  
*NASA TM—2012-104606*, **28**, 117 pp.

Moorthi, S., and M. J. Suarez, 1992: Relaxed Arakawa-Schubert: A parameterization of moist  
convection for general circulation models. *Mon. Wea. Rev.*, **120**, 978-1002.

Murakami, H., B. Wang, and A. Kitoh, 2011: Future change of western north Pacific typhoons:  
Projections by a 20-km-mesh global atmospheric model. *J. Climate*, **24**, 1154-1169.

Noda, A. T., K. Oouchi, M. Satoh, and H. Tomita, 2012: Quantitative assessment of diurnal  
variation of tropical convection simulated by a global nonhydrostatic model without  
cumulus parameterization. *J. Climate*, **25**, 5119-5134.

Oouchi, K., J. Yoshimura, H. Yoshimura, R. Mizuta, S. Kusunoki, and A. Noda, 2006: Tropical  
cyclone climatology in a global-warming climate as simulated in a 20-km mesh global  
atmospheric model: Frequency and wind intensity analysis. *J. Meteor. Soc. Japan*, **84**, 259–  
276.

Putman, W. M., and M. Suarez, 2011: Cloud-system resolving simulations with the NASA  
Goddard Earth Observing System global atmospheric model (GEOS-5). *Geophys. Res. Lett.*,  
**38**, L16809, doi:10.1029/2011GL048438.

Rayner, N. A., and Coauthors, 2003: Global analyses of sea surface temperature, sea ice, and  
night time air temperature since the late nineteenth century. *J. Geophys. Res.*, **108**, (D14),  
4407, doi:10.1029/2002JD002670.

Rienecker M. M. and Coauthors, 2008: The GEOS-5 data assimilation system – documentation  
of versions 5.0.1 and 5.1.0, and 5.2.0. In: NASA technical report series on global modeling  
and assimilation, NASA/TM-2008-104606, **27**, 92pp.



Rienecker, M.M., and Coauthors, 2011: MERRA - NASA's Modern-Era Retrospective Analysis for Research and Applications. *J. Climate*, **24**, 3624-3648. doi: 10.1175/JCLI-D-11-00015.1.

Reed, K. A., and C. Jablonowski, 2011: Impact of physical parameterizations on idealized tropical cyclones in the Community Atmosphere Model. *Geophys. Res. Lett.*, **38**, L04805, doi:10.1029/2010GL046297.

Shaevitz, D. A., and Coauthors, 2014: Characteristics of tropical cyclones in high-resolution models in the present climate. Submitted to *J. Climate*,

Shen, B.-W., R. Atlas, O. Reale, S.-J. Lin, J.-D. Chen, J. Chang, C. Henze, and J.-L. Li, 2006: Hurricane forecasts with a global mesoscale-resolving model: Preliminary results with Hurricane Katrina (2005). *Geophys. Res. Lett.*, **33**, L13813, doi:10.1029/2006GL026143.

Simpson, J. , and V. Wiggert, 1969: Models of precipitating cumulus towers, *Mon. Wea. Rev.*, **97**, 471-480.

Slingo, J., and Coauthors, 1994: Mean climate and transience in the tropics of the UGAMP GCM: Sensitivity to convective parameterization. *Quart. J. Roy. Meteor. Soc.*, **120**, 881-922.

Smith, R. K., 2000: The role of cumulus convection in hurricanes and its representation in hurricane models. *Rev. Geophys.*, **38**, 465-489.

Stan, C., 2012: Is cumulus convection the concertmaster of tropical cyclone activity in the Atlantic? *Geophys. Res. Lett.*, **39**, doi: 10.1029/2012GL053449.

Stephens, G., and Coauthors, 2002: The CloudSat mission and the A-Train. *Bull. Amer. Meteor. Soc.*, **83**, 1771-1790.

Strachan, J., P.-L. Vidale, K. Hodges, M. Roberts, and M.-E. Demory, 2013: Investigating global tropical cyclone activity with a hierarchy of AGCMs: The role of model resolution. *J. Climate*, **26**, 133–152.

562 Tiedke, M., 1993: Representation of clouds in large-scale models. *Mon. Wea. Rev.*, **121**, 3040-  
563 3061.

564 Tokioka, T., K. Yamazaki, A. Kitoh, and T. Ose, 1988: The equatorial 30–60 day oscillation and  
565 the Arakawa-Schubert penetrative cumulus parameterization, *J. Meteorol. Soc. Japan.*, **66**,  
566 883–901.

567 Vitart, F., J. L. Anderson, J. Sirutis, and R. E. Tuleya, 2001: Sensitivity of tropical storms  
568 simulated by a general circulation model to changes in cumulus parameterization. *Quart. J.*  
569 *Royal Meteor. Soc.*, **127**, 25-51.

570 Vitart, F., D. Anderson, and T. Stockdale, 2003: Seasonal forecasting of tropical cyclone landfall  
571 over Mozambique. *J. Climate*, **16**, 3932-3945.

572 Walsh, K. J. E., M. Fiorino, C. W. Landsea, and K. L. McInnes, 2007: Objectively determined  
573 resolution-dependent threshold criteria for the detection of tropical cyclones in climate  
574 models and reanalyses. *J. Climate*, **20**, 2307–2314.

575 Walsh, K., S. Lavender, E. Scoccimarro, and H. Murakami, 2013: Resolution dependence of  
576 tropical cyclone formation in CMIP3 and finer resolution models. *Clim. Dyn.*, **40**, 585-599.

577 Wang, Z., 2014: Role of cumulus congestus in tropical cyclone formation in a high-resolution  
578 numerical model simulation. *J. Atmos. Sci.*, early online release.

579 Zhang, G. J., and N. A. McFarlane, 1995: Sensitivity of climate simulations to the  
580 parameterization of cumulus convection in the Canadian climate centre general circulation  
581 model. *Atmosphere-Ocean*, **33(3)**, 407-446.

582 Zhao, M., I. M. Held, S.-J. Lin, and G. A. Vecchi, 2009: Simulations of global hurricane  
583 climatology, interannual variability and response to global warming using a 50-km  
584 resolution GCM. *J. Climate*, **22**, 6653-6678.

585 Zhao, M., I. M. Held, and S.-J. Lin, 2012: Some counter-intuitive dependencies of tropical  
586 cyclone frequency on parameters in a GCM. *J. Atmos. Sci.*, **69**, 2272-2283.

587

588 **Table 1** The number of TCs (middle and bottom row) detected for each year and experiment.

589 Rightmost column represents the observed TC numbers archived at International Best Track

590 Archive for Climate Stewardship (IBTrACS) (Knapp et al. 2010). The values in parenthesis in

591 the middle and bottom row are the number of hurricanes.

	ExpA	ExpB	ExpC (control run)	Observation
2005	36 (24)	25 (15)	18 (8)	29 (15)
2006	27 (21)	13 (8)	8 (6)	10 (5)

592

593

**Table 2.** List of the threshold values for detecting the TC using the detection/tracking algorithm based on Vitart et al. (2003)

Variables	local relative vorticity maximum (850hPa)	warm core temperature maximum	minimum sea level pressure (SLP)	minimum lower-level wind speed	minimum duration
Criteria	$1.0 \times 10^{-4} \text{ s}^{-1}$	6K, Distance between the TC center and the center of the warm core must not exceed $2^\circ$ lon.&lat.	Minimum SLP defines the TC center and must exist within $2^\circ \times 2^\circ$ radius of the vorticity maximum	$17 \text{ m s}^{-1}$	4 days

602

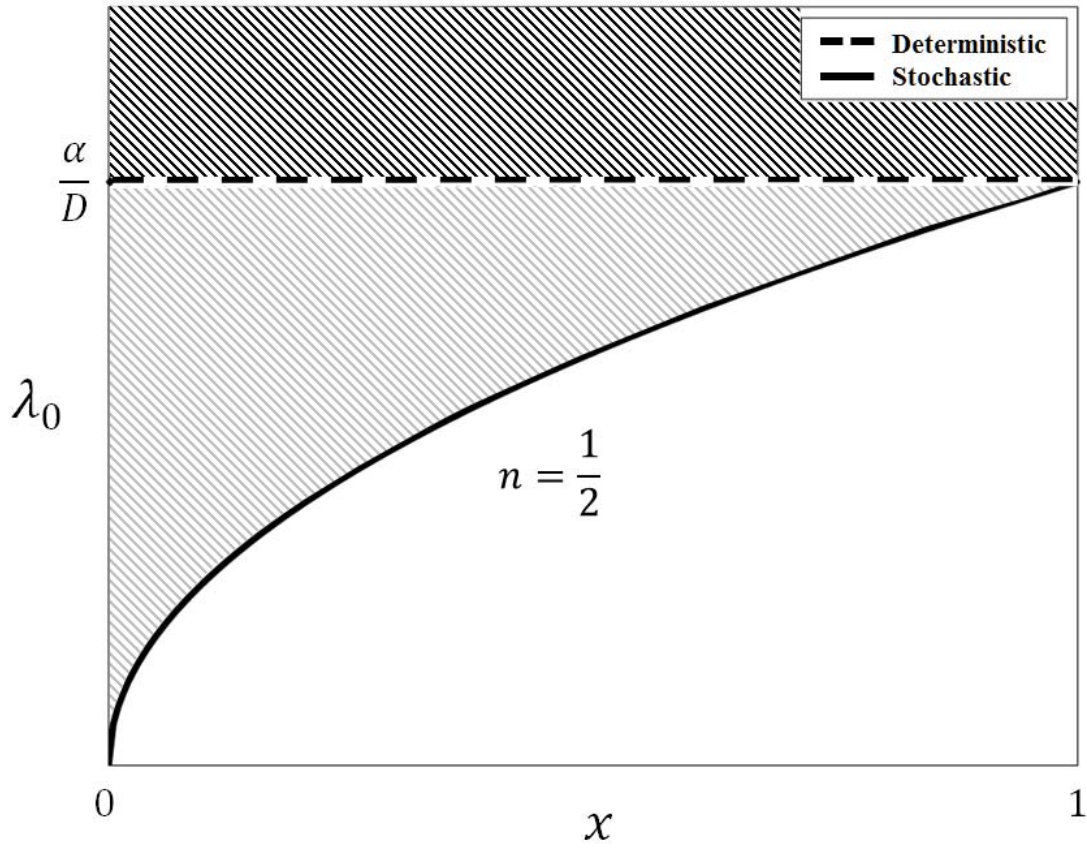
603 **Table 3.** Difference in total precipitation averaged over the North Atlantic basin that covers  
 604 90°W-0°E and 0°N-50°N. Difference between ExpA and control run is shown in the middle  
 605 column whereas the difference between ExpB and control run in the right column. The second  
 606 and the third row represent the result for the year 2005 and 2006, respectively. Observed  
 607 precipitation on the left column is obtained from the Global Precipitation Climatology Project  
 608 (GPCP) dataset, which is 2.5 lon.-lat. gridded.

	ExpA minus ExpC	ExpB minus ExpC
2005 (Obs.=3.62)	4.89-4.40=0.49	4.43-4.40=0.03
2006 (Obs.=3.57)	5.02-4.42=0.60	4.51-4.42=0.09

609

610

611



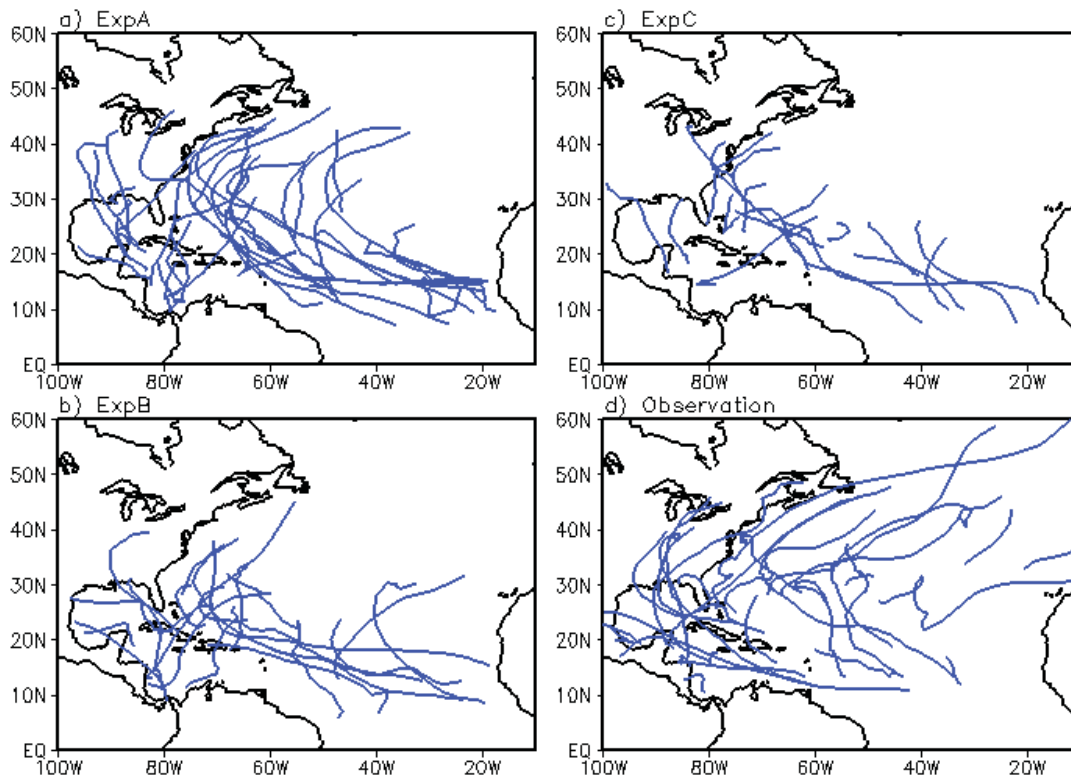
612

613

614 Figure 1 Profile of  $\lambda_0$  as a function of  $x$  in equation (1) with  $n=1/2$  (solid curve). A straight  
 615 dashed line represents the deterministic  $\lambda_0=\alpha/D$ , which applies the Tokioka constraint identically  
 616 at every time step. The entrainment rate of convective plumes need to be higher than  $\lambda_0$ , which  
 617 range is represented with the black hachured area for the deterministic method. The entrainment  
 618 rate allowed for the stochastic method is represented with both the black and grey hachured areas.

619

Atlantic TSC tracks (2005)



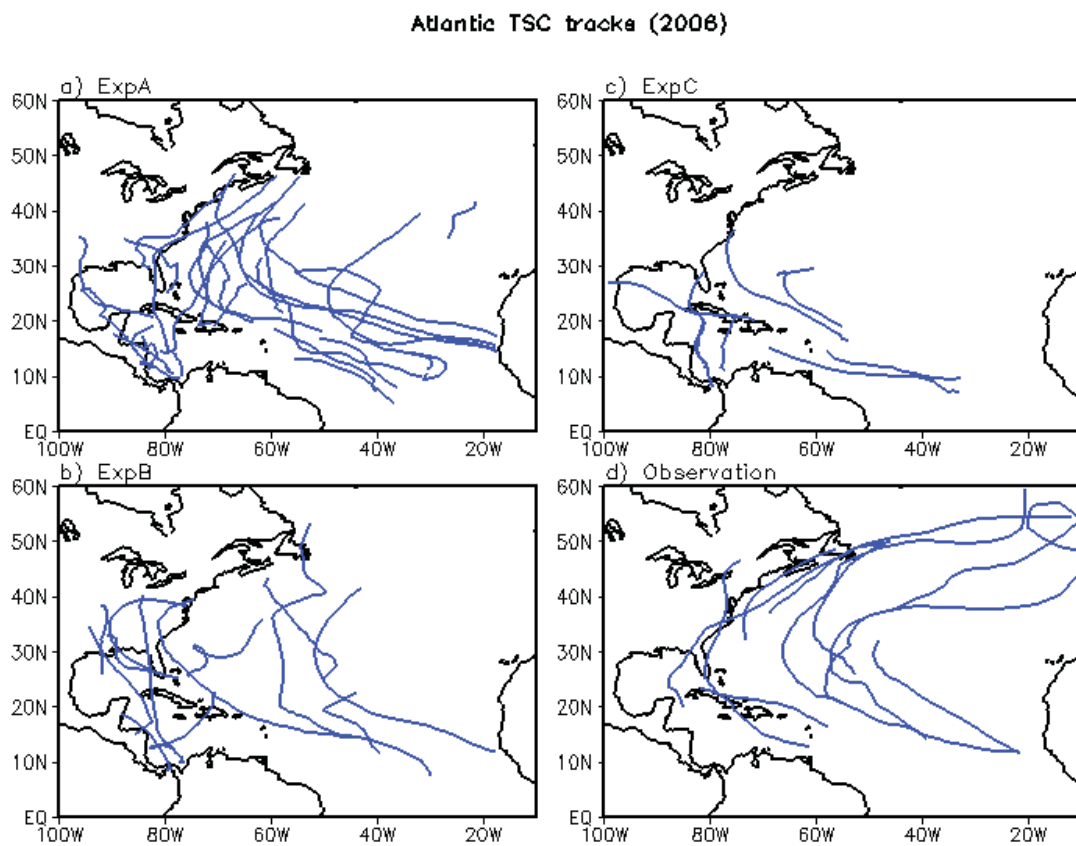
620

621

622 Figure 2. TC tracks for 2005 for a) ExpA, b) ExpB, c) ExpC and d) Observations.



623



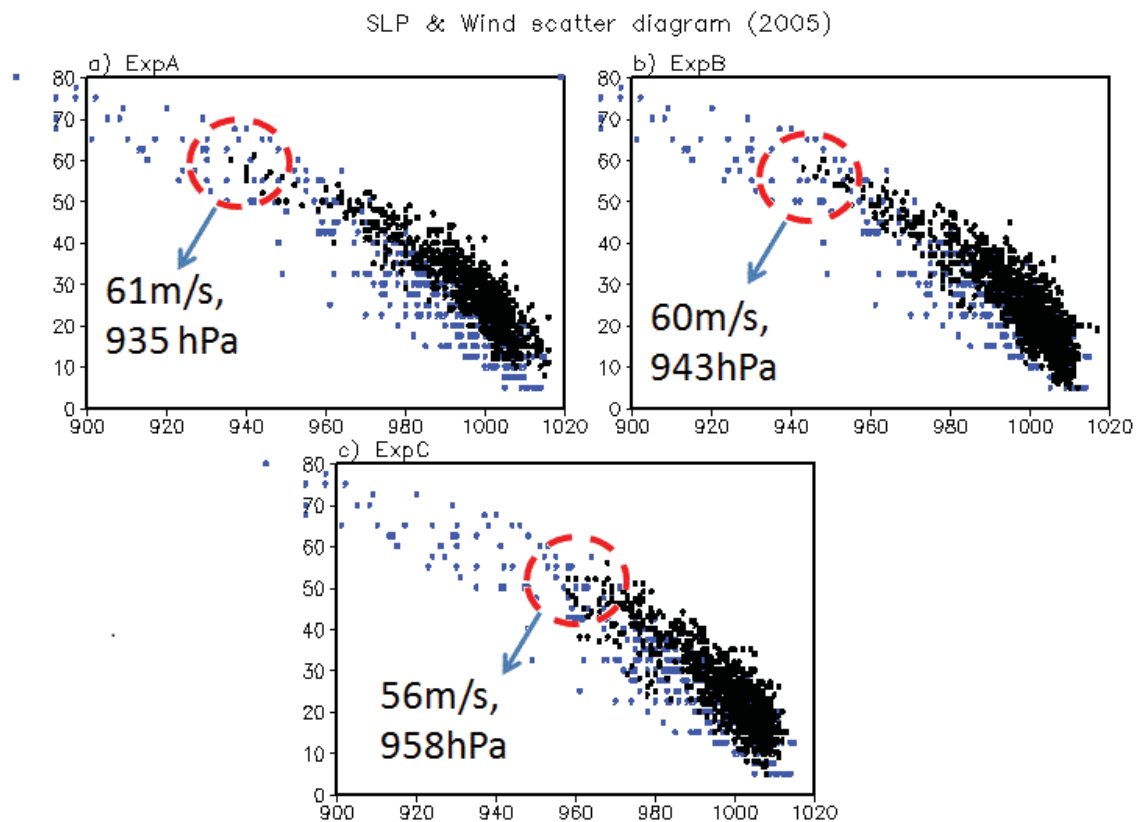
624

625

626 Figure 3. Same as Figure 2, but for 2006.

627

628



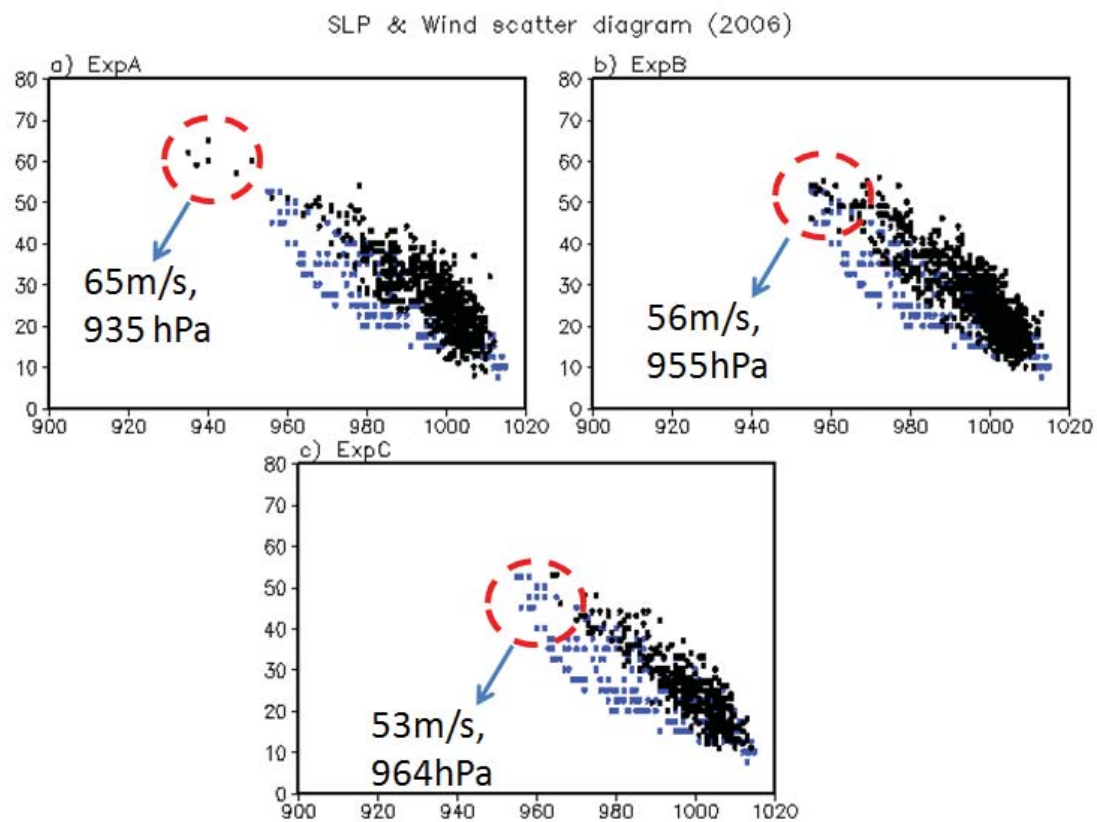
629

630

631 Figure 4. Scatter distribution of sea level pressure (SLP) and lower-level wind speed of TCs in  
632 2005 for a) ExpA, b) ExpB and c) ExpC.

633

634



635

636

637 Figure 5. Same as Figure 4 but for 2006.

638

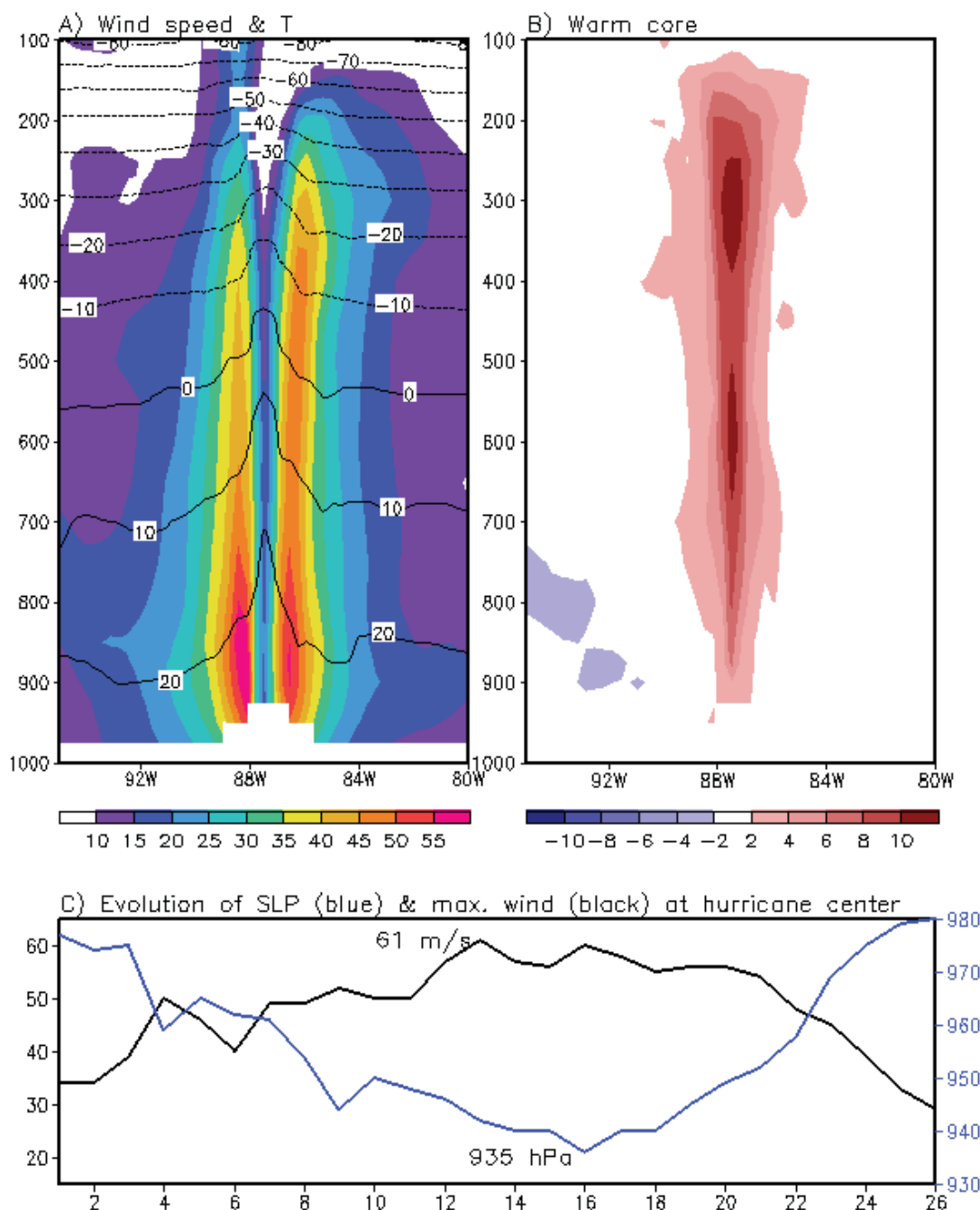


Figure 6. a) and b): Longitude-height cross-section of the strongest hurricane in 2005 simulated by ExpA. Shaded and contoured in a) is wind speed and temperature, respectively, and b) represents warm core computed as temperature deviation from zonal mean over 40° longitudes with a hurricane core centered. Black line on the bottom panel denotes the time evolution of the

maximum wind speed at the lower-level (850-950hPa) of hurricane center whereas blue line the minimum SLP.

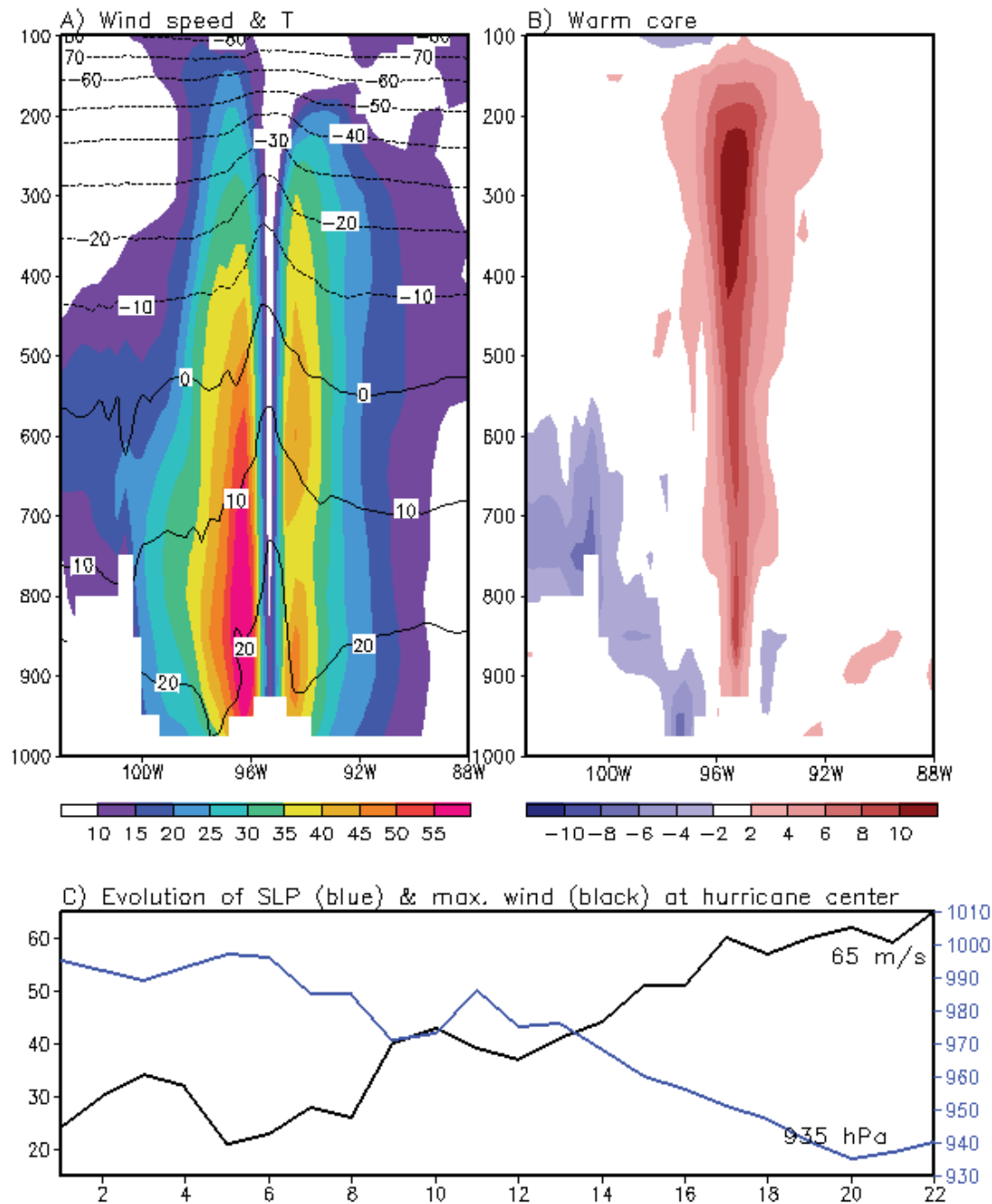


Figure 6: (Continued) Same as the figure on the previous page but for 2006.

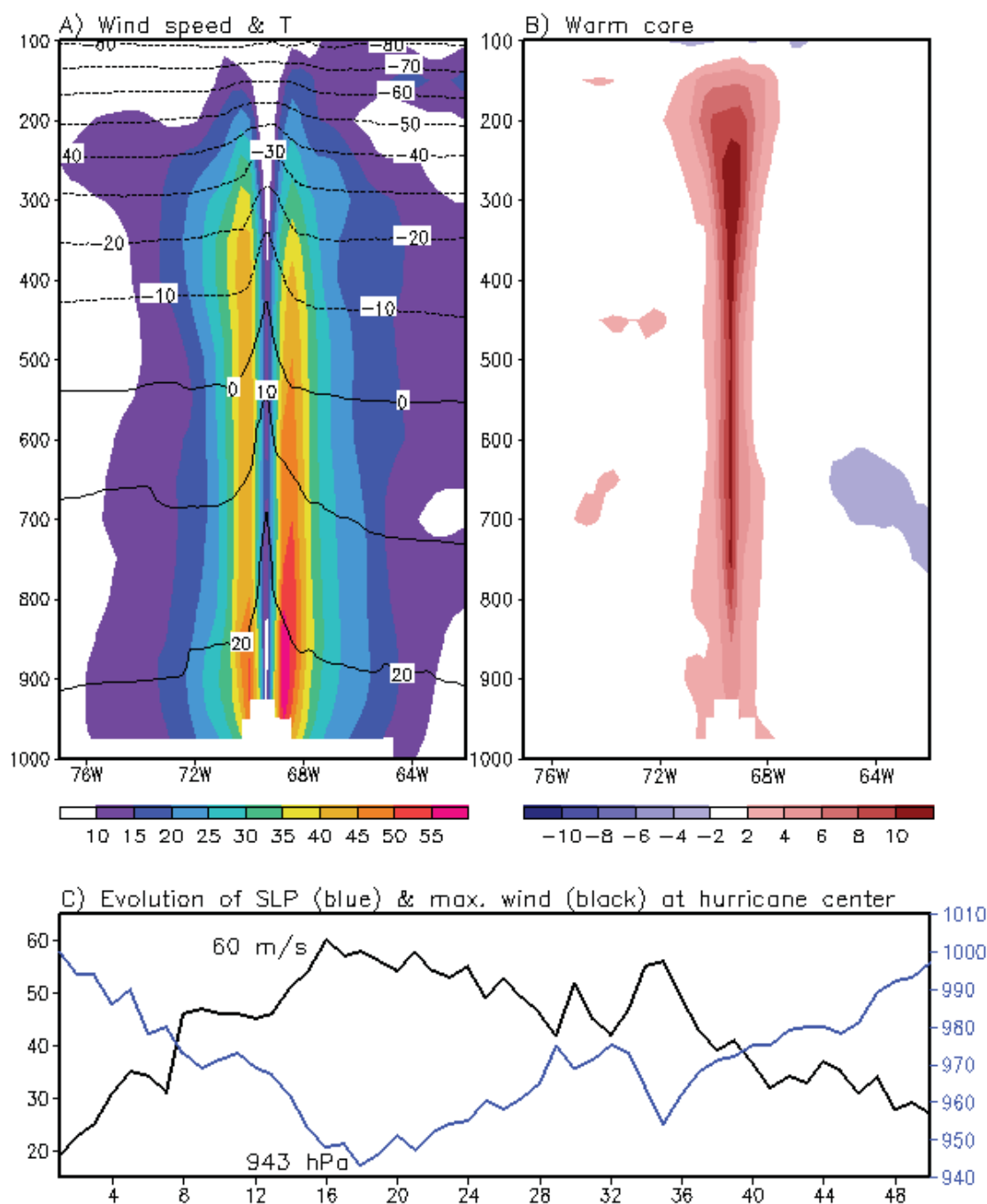


Figure 7. Same as Figure 6 but for ExpB

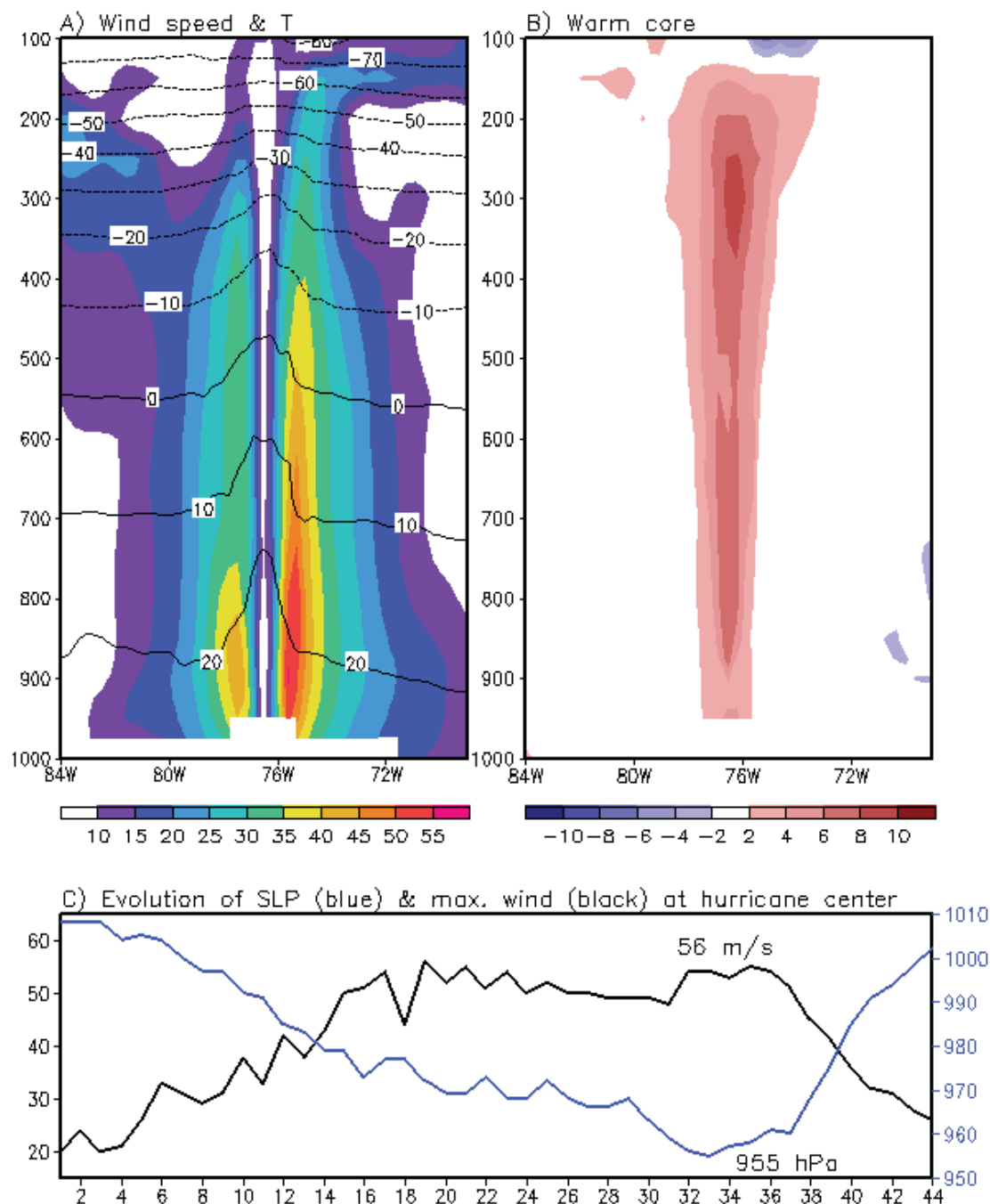


Figure 7: (Continued) Same as the figure on the previous page but for 2006.

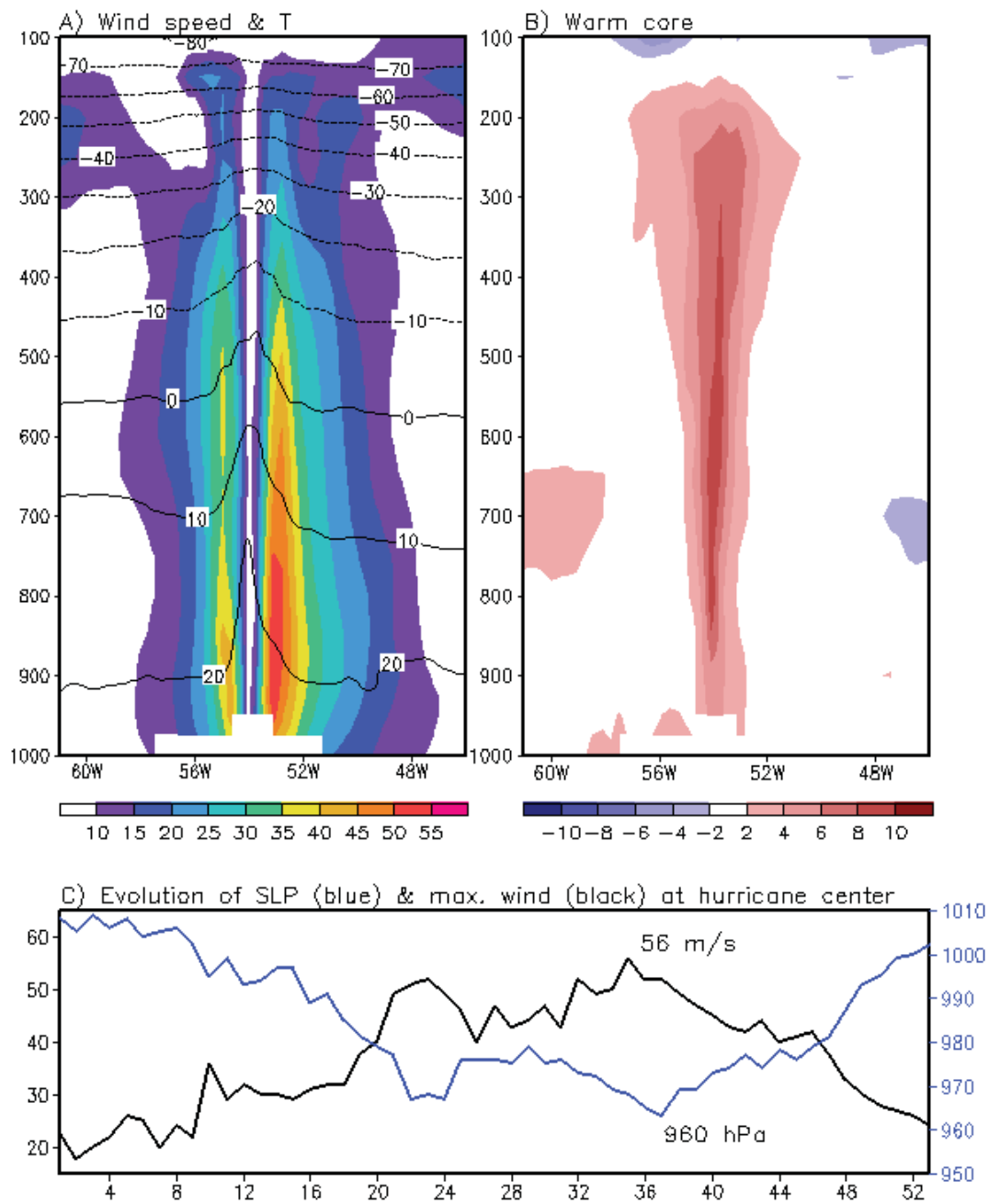


Figure 8. Same as Figure 6 but for ExpC.



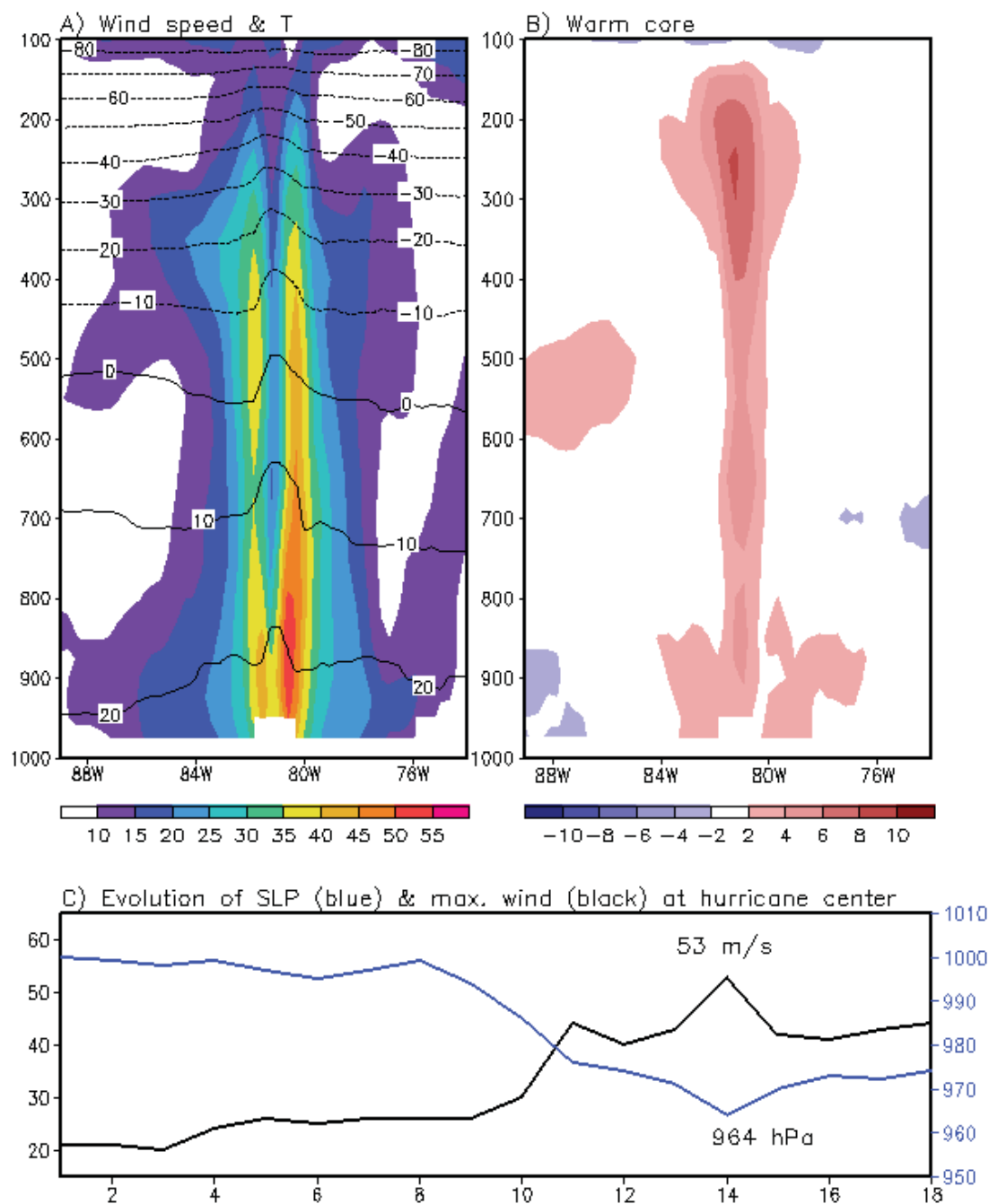


Figure 8: (Continued) Same as the figure on the previous page but for 2006.

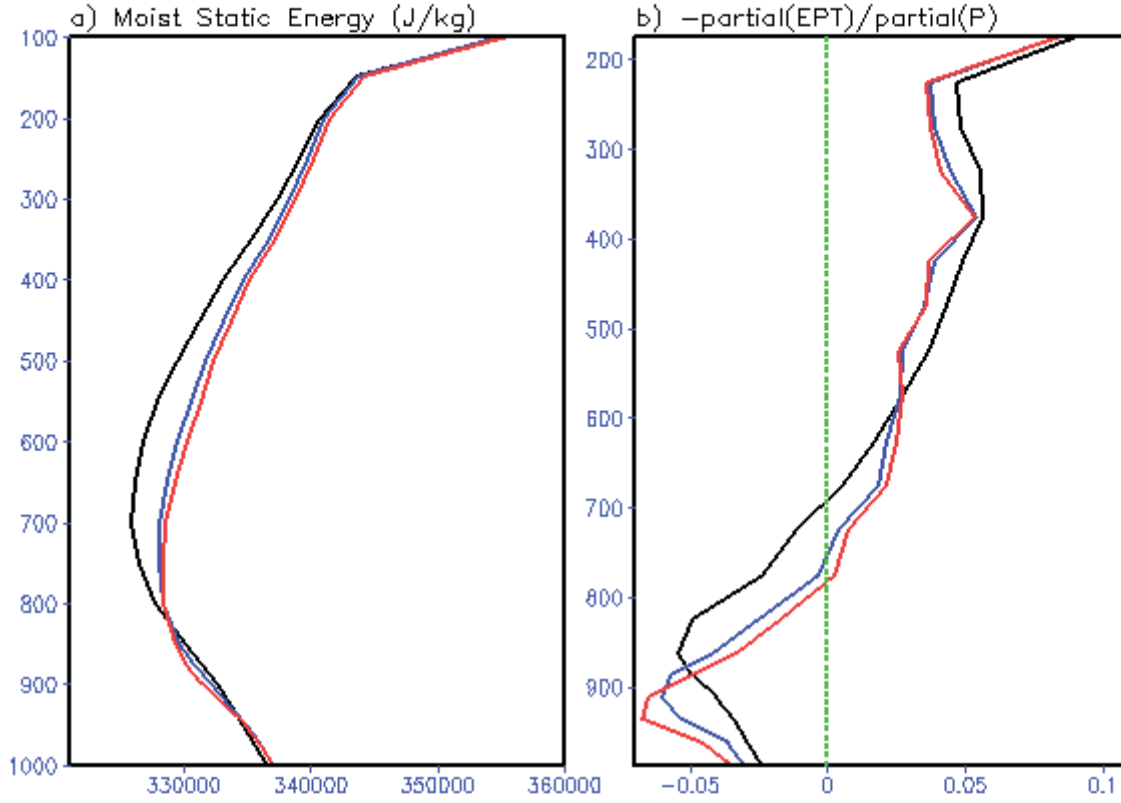


Figure 9. Vertical profile of area-averaged a) moist static energy (J/Kg) and b) vertical gradient of equivalent potential temperature ( $-\frac{\partial \theta_e}{\partial P}$ ) from ExpA (black), ExpB (blue) and ExpC (red) experiments for 2005. Geographical domain for area-averaging covers the typical TC genesis region over the Atlantic that spans 60°W-15°W, 5°N-20°N.

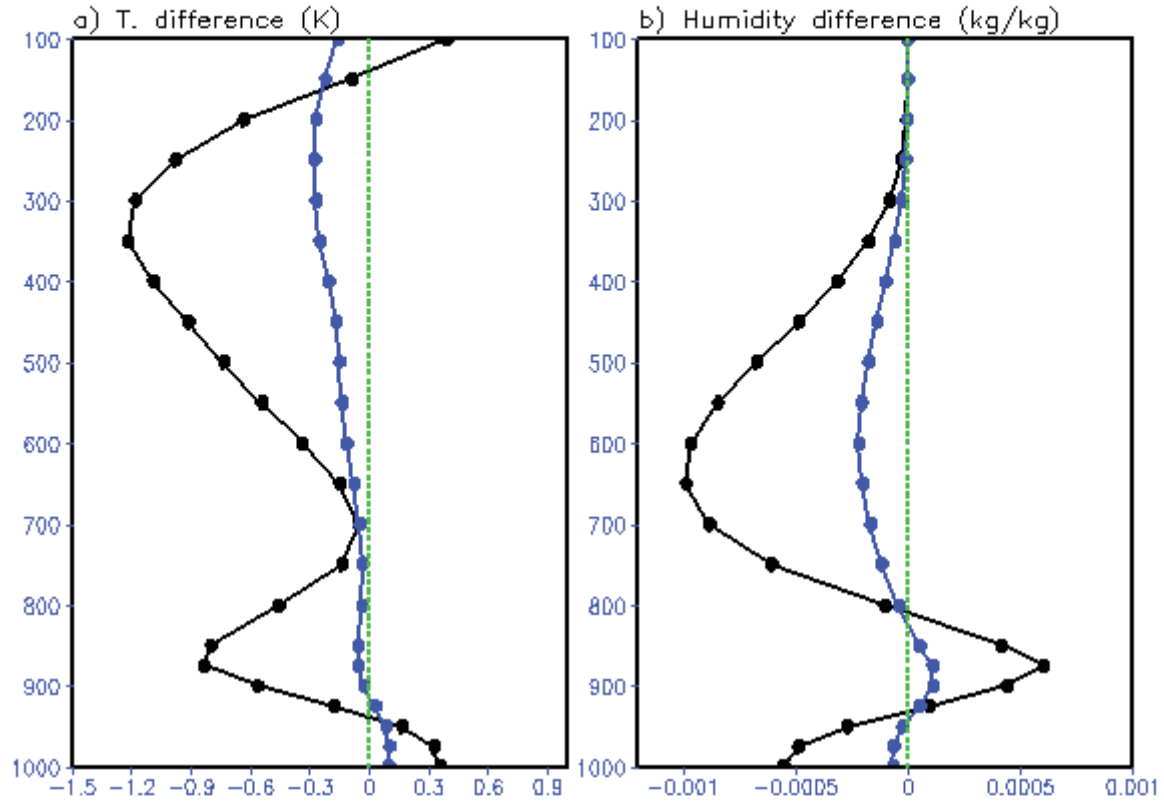


Figure 10. Left: Vertical profile of area-averaged difference in a) temperature and b) specific humidity for 2005. Black curve denotes difference between ExpA and ExpC (ExpA minus ExpC) while the blue curve denotes difference between ExpB and ExpC (ExpB minus ExpC). Geographical domain for area-averaging covers the typical TC genesis region over the Atlantic that spans 60°W-15°W, 5°N-20°N.

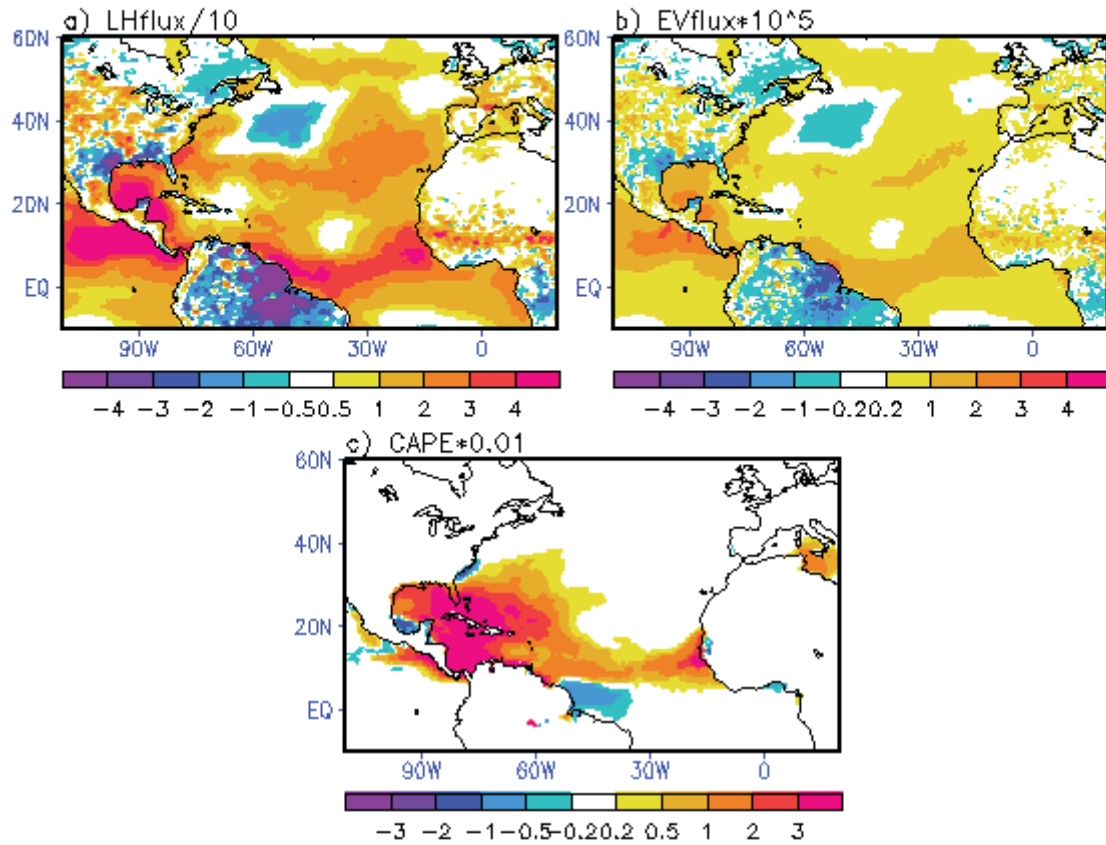


Figure 11. Difference in a) surface latent heat flux, b) evaporative flux, and c) convective available potential energy between ExpA and ExpC (ExpA minus ExpC) during hurricane season (June through November) of 2005.

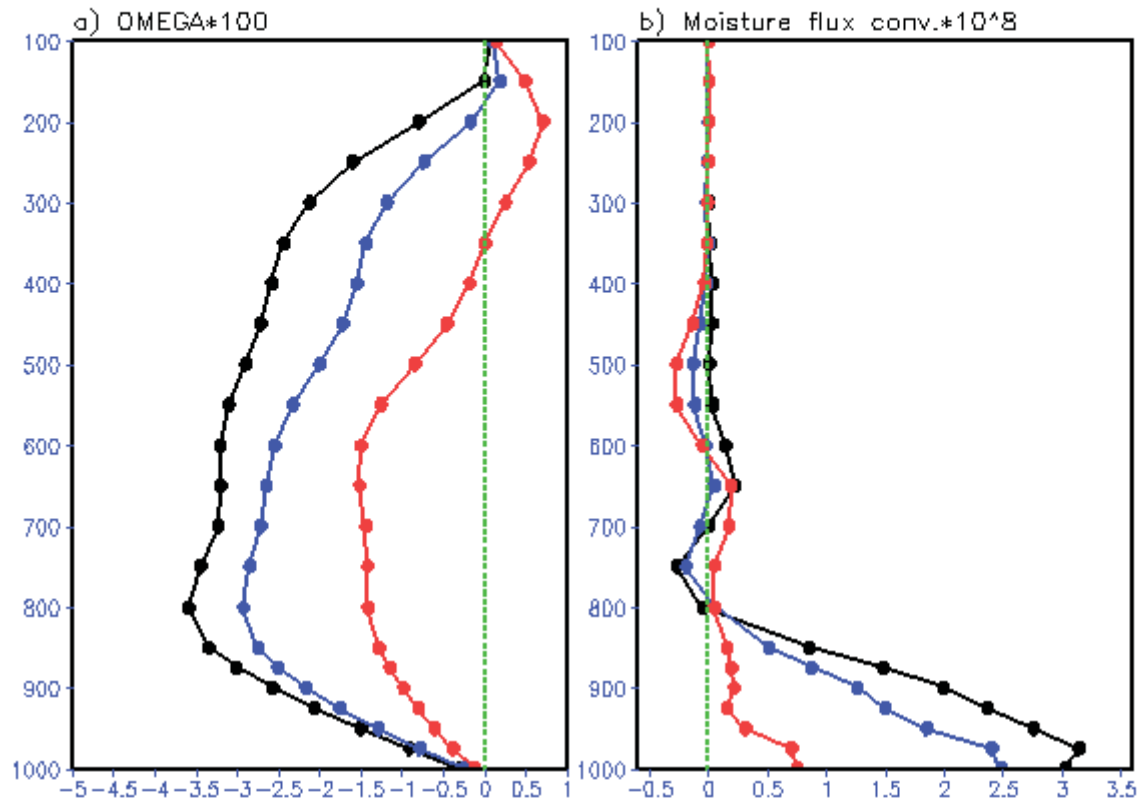


Figure 12. The vertical profile of area-averaged a) omega velocity and b) moisture flux convergence for the hurricane season (June through November) of 2005. The geographical domain for area-averaging covers the typical TC genesis region over the Atlantic that spans 60°W-15°W, 5°N-20°N. Black, blue and red curves denote results for ExpA, ExpB and ExpC, respectively.

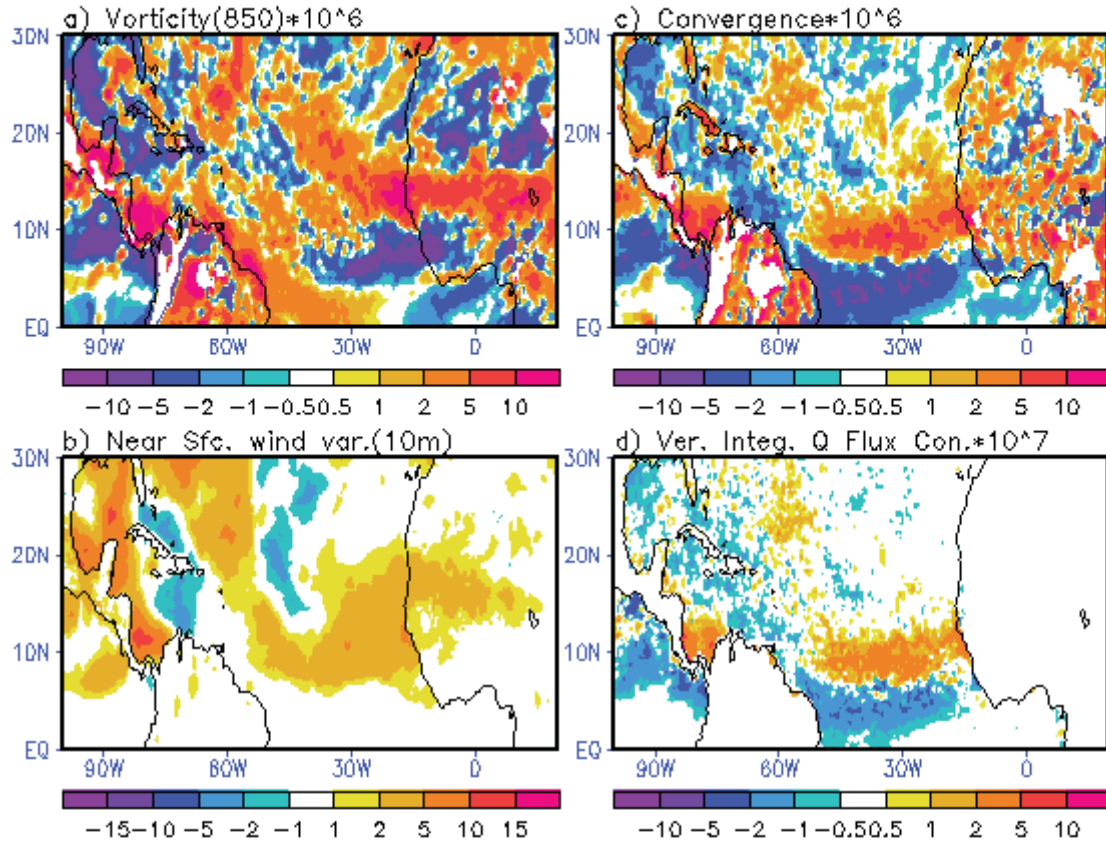


Figure 13. Difference in a) relative vorticity at 850hPa, b) wind variance at 10m level, c) convergence at 925hPa and d) vertically integrated moisture flux convergence between ExpA and ExpC (ExpA minus ExpC) during hurricane season (June through November) of 2005.

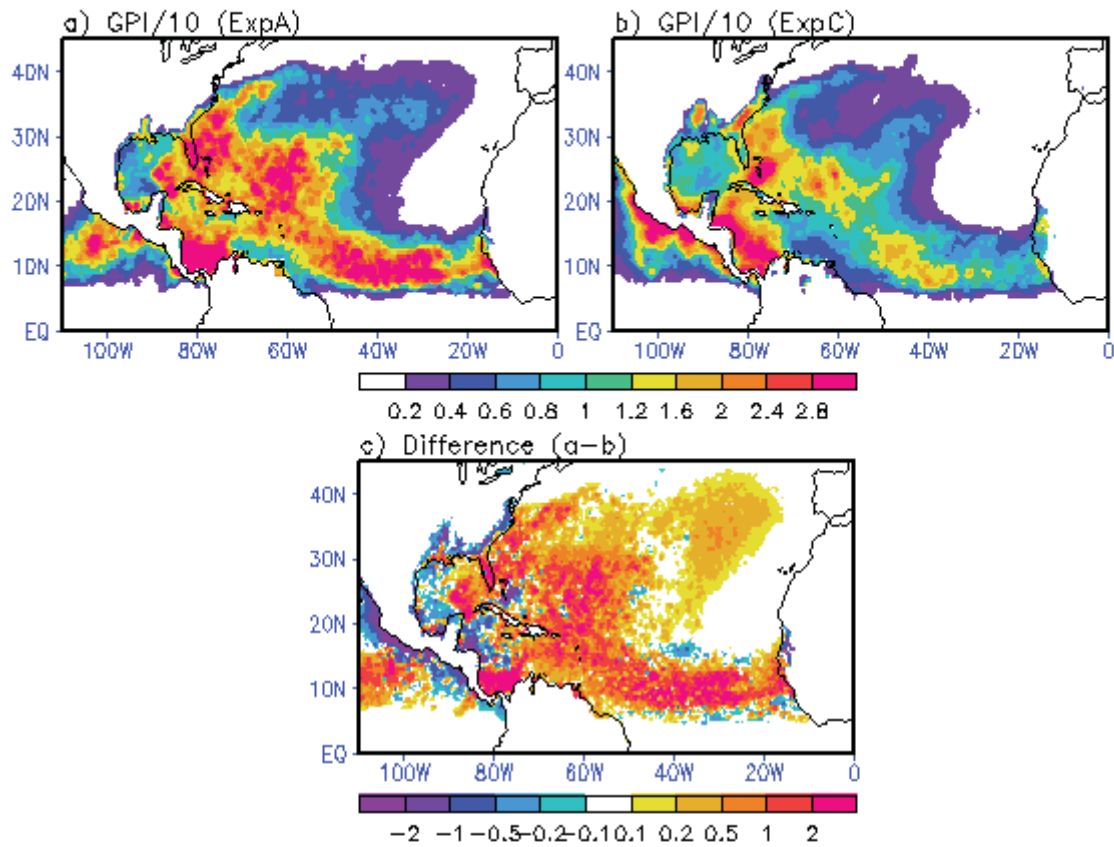


Figure 14. Distribution of modified genesis potential indices for 2005 from a) ExpA, b) ExpC and c) difference (ExpA minus ExpC)

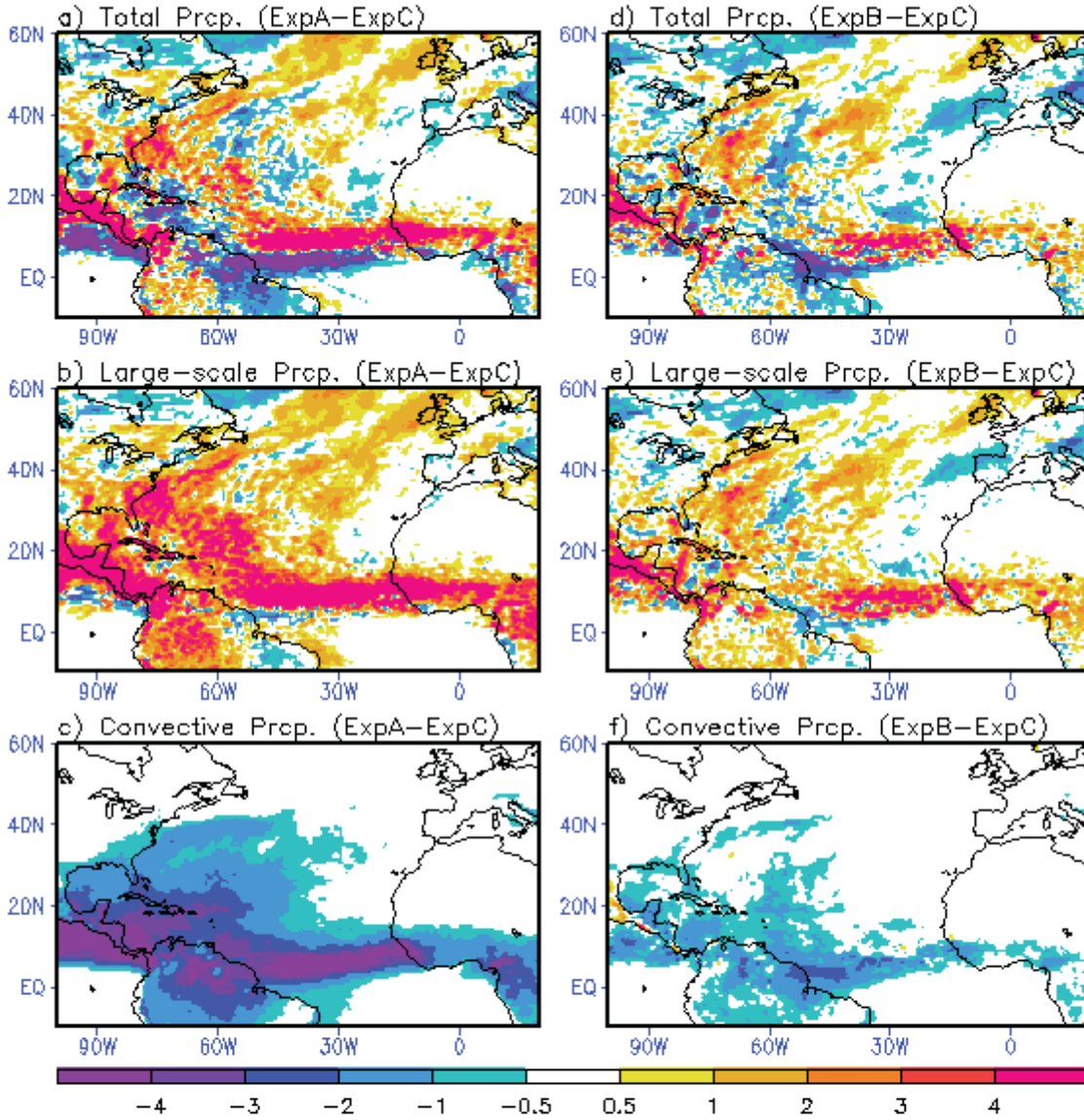


Figure 15. Difference in a) total precipitation, b) large-scale precipitation and c) convective precipitation between ExpA and ExpC (ExpA minus ExpC) during 2005 hurricane season. Right panel is the same as the left panel but for difference between ExpB and ExpC (ExpB minus ExpC).



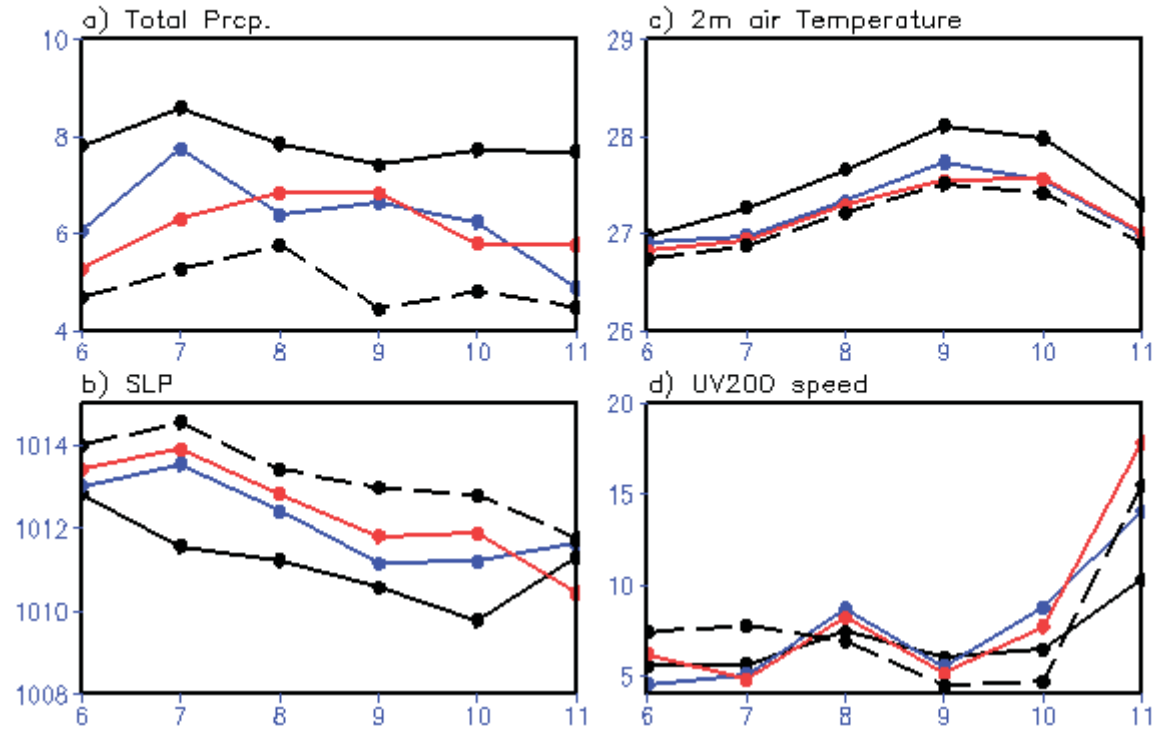


Figure 16. The seasonal variation of the area-averaged a) total precipitation, b) SLP, c) 2m air temperature and d) 200hPa wind speed for 2005. The geographical domain for area-averaging covers the typical TC genesis region over the Atlantic that spans 60°W-15°W, 5°N-20°N. Black, blue, and red solid lines represent the results from ExpA, ExpB and ExpC, respectively. The observed (GPCP) precipitation and remaining variables from MERRA are plotted with black dashed lines.